Spray-Deposited Titanium-Oxide Films For Infrared Optics, Photonics, And Solar Cell Applications.

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SPRAY-DEPOSITED TITANIUM-OXIDE FILMS FOR INFRARED OPTICS, PHOTONICS, AND SOLAR CELL APPLICATIONS

by

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ABSTRACT

Self-assembled TiO$_2$ foam-like films, were grown by the water based Streaming Process for Electrodeless Electrochemical Deposition (SPEED). The morphology of the $\sim 1 \mu$m thick films consists of a tangled ropy structure with individual strands of $\sim 200$ nm diameter and open pores of 0.1 to 3 micron dimensions. Such films are advantageous for proposed perovskite solar cell comprising CH$_3$NH$_3$PbI$_3$ absorber with additional inorganic films as contact and conduction layers, all deposited by SPEED. Lateral film resistivity is in the range 20 - 200 k$\Omega$-cm, increasing with growth temperature, while sheet resistance is in the range $2 - 2 \times 10^8 \Omega$/Sq. Xray diffraction confirms presence of TiO$_2$ crystals of orthorhombic class (Brookite). UV-vis spectroscopy shows high transmission below the expected 3.2 eV TiO$_2$ bandgap. Transmittance increases with growth temperature. This is a Ropy TiO$_2$ thin film.

We also prepared a Smooth TiO$_2$ thin film. Self-assembled TiO$_2$ film deposited by aqueous-spray deposition was investigated to evaluate morphology, crystalline phase, and infrared optical constants. The Anatase nano-crystalline film had $\sim 10$ nm characteristic surface roughness sparsely punctuated by defects of not more than 200 nm amplitude. The film is highly transparent throughout the visible to wavelengths of 12 $\mu$m. The indirect band gap was determined to be 3.2 eV. Important for long-wave infrared applications is that dispersion in this region is weak compared with the more commonly used dielectric SiO$_2$ for planar structures. The low-cost, large-area, atmospheric-pressure, chemical spray deposition method allows conformal fabrication on flexible substrates for long-wave infrared photonics.

For comparison TiO$_2$ films deposited by electron-beam evaporation were evaluated to determine morphology, crystalline phase, and optical transparency.
The evaporated TiO$_2$ film was amorphous but crystallized into Anatase phase after annealing.

Such film is attractive as electron conductor of unprecedented thinness and flexibility for proposed perovskite solar cell comprising CH$_3$NH$_3$PbI$_3$ absorber with additional inorganic films as contact and conduction layers. The spray deposition method would allow conformal solar cell fabrication on flexible substrates for wearable power generation. Band gap of Evaporated TiO$_2$ film is 4.0 eV. We prepared BaTiO$_3$ thin film to know infrared pyroelectric response.

Self-assembled nano-crystalline BaTiO$_3$ films on stainless steel foil substrates, were grown by the water based Streaming Process for Electrodeless Electrochemical Deposition (SPEED). SPEED is an aqueous process that deposits self-assembled nano-crystalline inorganic thin films over large areas, without a vacuum, providing a scalable and manufacturing friendly process to fabricate durable films. The morphology of the $\sim$1$\mu$m thick films comprises single crystals of micron dimensions imbedded in a matrix of nanocrystals. XRD confirms presence of BaTiO$_3$ crystals of hexagonal phase for samples annealed at 500°C. Subsequent annealing at 600°C transforms the film to the cubic phase. Potential applications include dielectric layers, capacitors, waveguides, ferroelectric RAM, pyroelectric infrared detectors, and phosphors. Characterization of infrared pyroelectric response at 10$\mu$m wavelength shows an initially good sensitivity that reversibly decays over a period of days due to water vapor absorption. A short-lived photo-response due to poling of the hydrated sample is also observed. We studied BaTiO$_3$ to know hysteresis loop.

Pyroelectric photoresponse of aqueous spray deposited thin films containing BaTiO$_3$ nano-crystals is reported. X-ray diffraction data indicate the presence of hexagonal BaTiO$_3$ nano-crystals with $\sim$20 nm crystalline domains in a matrix of some as yet unidentified nano-crystalline material. When the film is annealed at 600°C, the X-ray pattern changes significantly and indicates a conversion to one of the non-hexagonal phases of BaTiO$_3$ as well as a complete change in the matrix. With suitable amplifier, the measured photoresponse was 40V/W.
Ferroelectric hysteresis on a film with significant presence of hexagonal BaTiO$_3$ shows saturated polarization which is about 5-times smaller than for the bulk tetragonal phase.
It is our genuine gratefulness and warmest regard that we dedicate this work to my country IRAQ (Mesopotamia), Higher Committee for Education Development in Iraq (HCED), my father, my mother, my wife, my kids, my brothers, my sister, who are close to me, and close to my family.
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# TABLE OF CONTENTS

LIST OF FIGURES .................................................. xiii

LIST OF ABBREVIATIONS ............................................ xvi

CHAPTER 1: INTRODUCTION .......................................... 1

CHAPTER 2: EXPERIMENTAL METHODS ............................. 5

2.1 Experimental Methods for Ropy TiO$_2$ for Solar Cells ......... 5

2.2 Experimental Methods for Smooth TiO$_2$ for Infrared Total Photonics ............................................ 8

2.3 Experimental Methods for SPEED BaTiO$_3$ for Pyroelectrics ............................................ 10

CHAPTER 3: RESULTS AND DISCUSSION OF ROPY TiO$_2$ FOR SOLAR CELLS ........................................ 15

CHAPTER 4: RESULTS AND DISCUSSION OF SMOOTH TiO$_2$ FOR INFRARED TOTAL PHOTONICS ............................... 19

CHAPTER 5: RESULTS AND DISCUSSION OF EVAPORATED TiO$_2$ ........................................ 30

CHAPTER 6: RESULTS AND DISCUSSION OF SPEED BaTiO$_3$ FOR PYROELECTRICS ........................................ 36
LIST OF FIGURES

2.1 Figure 2.1-a is a schematic of proposed Perovskite Solar Cell layer structure. Figure 2.1-b is an optical microscope image of SPEED-grown TiO$_2$ thin film. .......................................................... 6

2.2 SPEED Deposition method .......................................................... 8

2.3 Photograph of experimental set up for photoresponse measurements. A spring contact connects to the carbon-black coating on the BaTiO$_3$ thin-film. The amplifier is constructed on the back side of the printed-circuit board. The arrow points to the spot of a red alignment laser which indicates the position illuminated by the pulsed CO$_2$ laser ........................................ 11

3.1 SEM images for TiO$_2$:Ni film. Figure 3.1-a is a top view. Figure 3.1-b is a cross section. .......................................................... 15

3.2 Figure 3.2-a is a typical XRD spectrum of SPEED-grown with Brookite phase reference data. Figure 3.2-b is a resistivity, sheet resistance, film thickness, and optical transmittance at 500 nm wavelength versus growth temperature. .......................................................... 17

4.1 XRD Smooth TiO$_2$ film and reference spectra .......................................................... 20

4.2 SEM image of Smooth SPEED TiO$_2$ film surface .......................................................... 21

4.3 SEM cross-sectional image of Smooth SPEED TiO$_2$ film. .................................................. 22
4.4 AFM images of Smooth SPEED TiO$_2$. .................................................. 23

4.5 Transmission spectrum of Smooth SPEED TiO$_2$. ................................. 24

4.6 Determination of band gap of Smooth SPEED TiO$_2$ film. ....................... 26

4.7 Infrared complex index spectra for Smooth SPEED TiO$_2$ and evaporated SiO$_2$ films. The shaded region indicates LWIR wavelengths. ............ 28

5.1 XRD measurement of Evaporated TiO$_2$ on Glass. ................................. 30

5.2 Band gap of E-beam evaporation. ............................................................ 31

5.3 Transmission of E-beam evaporation ........................................................ 32

5.4 SEM Top View of Evaporated TiO$_2$. ....................................................... 33

5.5 SEM cross-sectional image of Evaporated TiO$_2$ film. ............................ 34

5.6 AFM 3-Dimensions of Evaporated TiO$_2$ Sample. ................................... 35

6.1 SEM images of SPEED-grown BaTiO$_3$ films. Figure 6.1-a is a top view revealing a mix of nano- and micro-meter scale crystals. Figure 6.1-b is a SEM cross section. ................................................................. 37

6.2 XRD spectra of SPEED-grown BaTiO$_3$ thin films.(left) BaTiO$_3$, 500°C, 2 hours annealed. (right) BaTiO$_3$, 600°C, 1 hour annealed. ................. 38
6.3 Figure 6.3-a is an oscilloscope screen shot showing the transient pyroelectric response to the laser pulse on channel 1. Channel 2 shows the gate of a boxcar average. Figure 6.3-b is a peak signal as function of position on the absorbing area of the BaTiO$_3$ device.

6.4 Figure 6.4-a is decay in peak photo-response with time for sample kept in air. Figure 6.4-b is a recovery of signal on baking at 100$^\circ$C.

6.5 High resolution SEM image of SPEED-grown BaTiO$_3$ thin film.

6.6 SEM image of surface of SPEED-grown BaTiO$_3$ film with carbon-black coating (lower region).

6.7 Reference Powder XRD spectra for the five phases of BaTiO$_3$.

6.8 XRD spectra of SPEED-grown BaTiO$_3$ thin films. Annealing temperatures in degrees C and the reference positions of XRD peaks for cubic and hexagonal phases are indicated.

6.9 Figure 6.9-a is an oscilloscope screen shot showing pyroelectric response to 100 mW CO$_2$ laser chopped at 70 Hz, amplifier A. Figure 6.9-b is a response to 13 mW/cm$^2$ broad-band IR source chopped at 1 Hz, amplifier B.

6.10 Ferroelectric hysteresis curve of BaTiO$_3$ thin film.
LIST OF ABBREVIATIONS

AFM .......... Atomic Force microscopy
ALD .......... Atomic Layer Deposition
CD .......... Conduction Band
CNC .......... Conduction Computer numerical control
EDS .......... Energy Dispersive Spectroscopy
ESL .......... Electron selective layers
FIB .......... Focus Ion Beam
HSL .......... Hole selective layers
LED .......... light-Emitting Diode
LWIR .......... Long-Wave Infrared
ICDD .......... International Centre for Diffraction Data
MCF .......... Materials Characterization Facility
PSC .......... Perovskite Solar Cells
PCE .......... Power conversion efficiencies
RAM .......... Random-Access Memory
SEM .......... Scanning electron microscopy
SPEED .......... Streaming Process for Electrodeless Electrochemical Deposition
VPSPEED ...... Vapor Phase Streaming Process for Electrodeless Electrochemical Deposition
CHAPTER 1: INTRODUCTION

Performance of Perovskite Solar Cells (PSC) has advanced rapidly[1, 2, 3, 4, 5, 6]. Perovskites are compounds of the form ABO$_3$. Solar Cells based on spin-coated mesoporous TiO$_2$ scaffold, CH$_3$NH$_3$PbI$_3$ perovskite absorber layer, and a 2, 2, 7, 7- tetrakis-(N, N-di-p-methoxyphenylamine)-9, 9-bifluorene (spiro-OMeTAD) hole transport layer have been reported with Power Conversion Efficiency (PCE) as high as 15%, and with 20% considered achievable. However, mesoporous TiO$_2$ as a conducting layer has demonstrated poor quality and poor electron conduction [7]. We present here a robust foam-like TiO$_2$ film grown by a water-based process that can be a superior alternative. Similarly, the organic spiro-OMeTAD used previously as hole conductor has poor hole transport. Instead, durable films of CuS with excellent hole transport grown by our same aqueous process might be substituted [8, 9].

Titanium dioxide is a very well-studied material with wide ranging applications, including electron collector in solar cells [10, 11, 12, 13], optical waveguides [14], optical coatings [15], polarized [16], ultra-thin capacitors [17], corrosion-barriers [18], protective coatings [19], MOSFETs [20], humidity sensors [21], and oxygen monitors [22]. Enabling features of TiO$_2$ include excellent visible transmittance, chemical stability, high permittivity [23], mechanical durability, high electrical resistance [24], and rigidity [25]. This work considers a Smooth TiO$_2$ as a potential useful material for long-wave infrared (LWIR) photonics applications. This consideration is motivated by the relatively low dispersion for TiO$_2$ in the 8 to 12$\mu$m wavelength range, in comparison to SiO$_2$, which is the most common dielectric in devices fabricated by silicon processing.
An additional motivation for this work is to evaluate a potentially advantageous aqueous deposition method that provides smooth conformal coatings. Titanium dioxide TiO$_2$ exists in nature as the minerals Rutile (tetragonal), Anatase (tetragonal) and Brookite (orthorhombic). Many factors affect the crystalline phase of deposited TiO$_2$ thin films, including deposition method, deposition temperature, annealing temperature, deposition rate, deposition pressure, precursor type, reaction atmosphere, impurities, and substrate type. TiO$_2$ films have been deposited by many different techniques, including hydrolysis [26], pyrolysis [27], pneumatic spraying [28], ultrasonic spraying [29], dip coating [30], plasma enhanced chemical vapor deposition [31], atmospheric pressure chemical vapor deposition [32], metal organic chemical vapor deposition [33], ultra high vacuum chemical vapor deposition [34], low pressure chemical vapor deposition [35], evaporation [36], spin on methods [37], sputtering [38], ion assisted deposition [38], plasma anodization [39], reactive ion plating [40], laser ablation [41], filtering arc deposition [25], atomic layer epitaxy [42], and screen printing [43]. The complex refractive index spectra of TiO$_2$ thin films prepared by atomic layer deposition (ALD) [44] and by reactive sputter deposition [45] have already been reported out to far infrared wavelengths (125 µm). Both films were prepared by vacuum-based physical deposition methods. The atmospheric-pressure aqueous inhomogeneous-chemical reaction-based deposition method investigated here has the potential for much lower cost, large area, perfectly conformal [46, 47] roll-to-roll manufacturing. Physical and optical properties of these films are investigated in this work with particular emphasis on application to infrared photonics, such as metal-insulator-metal plasmonic resonant absorbers [48] and planar waveguides [49].

The inorganic compound barium titanate (BaTiO$_3$) is a member of the Perovskite family, which has the general formula ABO$_3$. BaTiO$_3$ has excellent dielectric, ferroelectric and piezoelectric properties. BaTiO$_3$ was the first ferroelectric ceramic, and it is a good candidate for a diverse applications due to its high permittivity and low loss [50].
Applications include pyroelectric detectors [51], capacitors [52], waveguide modulators [53], gate dielectrics [54] and holographic memory [55].

The optimum synthesis method for BaTiO$_3$ depends on the end application, as the method used has a significant impact on the structure and properties of BaTiO$_3$. Its five known phases are hexagonal, cubic, tetragonal, orthorhombic, and rhombohedral. This work considers the pyroelectric response to infrared radiation of BaTiO$_3$ thin films.

Pyroelectric response is the generation of small voltage difference across opposite faces of the film due to a change in temperature that changes the films intrinsic polarization. Of its 5 possible crystal classes, only the cubic phase is incapable by symmetry of being pyroelectric [56]. Methods of preparing BaTiO$_3$ include solid-state reaction [57], chemical methods [58], hydrothermal growth [59], sol-gel processing [60], spray pyrolysis [61], the oxalate route [62], microwave heating [28], a micro-emulsion process [63] a polymeric precursor method [64], nanometer oxides doping [63] and modified combustion process [65]. We present a different method of thin film preparation, which is superficially similar to spray pyrolysis, except that it does not involve thermal splitting of molecules, and hence operates at lower temperature.

Motivation for the subject research is to develop an in-house IR detector material that can be deposited, patterned, and integrated with on-chip infrared photonic and plasmonic devices, such as spectral sensors [49, 66]. Barium titanate (BaTiO$_3$) was selected because this well-known ferroelectric material is amenable to our spray deposition technique, known as Streaming Process for Electrodeless Electrochemical Deposition (SPEED). SPEED is superficially similar to spray pyrolysis, except that it does not involve thermal splitting of molecules, and hence operates at lower temperature. SPEED results in perfectly conformal nano-crystalline thin films with strong adhesion [10, 67, 68]. Pyroelectric photoresponse of SPEED grown BaTiO$_3$ thin films was first reported in ref [69].
A purpose of the present work is to further investigate this response and the dependence of materials properties on processing conditions. SPEED deposits inorganic thin films over large areas, without a vacuum, a scalable and manufacturing friendly process.

Water-soluble compounds with complexing agents grow films by heterogeneous reaction on the substrate, with little wasteful homogeneous reaction. Hydrophilic substrates bind hydroxyl ions, which serve as attachment and nucleation sites with density exceeding $10^{12} \text{cm}^{-3}$. The nebulized aqueous precursor impinges on a substrate heated to $\sim 300^\circ \text{C}$, giving growth rate exceeding 200 nm/min. Substrate heating provides reaction activation energy, and it decomposes and volatilizes reaction byproducts. The five known phases of bulk crystalline BaTiO$_3$ are hexagonal, cubic, tetragonal, orthorhombic, and rhombohedral. Tetragonal is the stable phase at room temperature. The cubic phase exists above the $120^\circ \text{C}$ Curie Temperature. The high temperature hexagonal phase can be meta-stable at room temperature. The orthorhombic and rhombohedral phases occur at lower temperatures. All but the cubic phase are capable of a pyroelectric photoresponse [56], which is the appearance of a small voltage difference across opposite faces of the material due to a temperature-induced change in the films intrinsic polarization. The actual stable phase of BaTiO$_3$ film at room temperature can differ from that of crystal due to interaction with the substrate, and it can depend on processing history. We previously reported both cubic and hexagonal phase for room temperature SPEED-grown BaTiO$_3$ thin films [69].
CHAPTER 2: EXPERIMENTAL METHODS

2.1 Experimental Methods for Ropy TiO$_2$ for Solar Cells

Our aqueous process is called Streaming Process for Electrodeless Electrochemical Deposition, or SPEED. This method deposits self-assembled nanocrystalline inorganic thin films over large areas without vacuum. Water-soluble compounds with complexing agents grow films by heterogeneous reaction on the substrate, with little wasteful homogeneous reaction. Hydrophilic substrates bind hydroxyl ions ($OH^-$), which are attachment sites for nucleation with density exceeding $10^{12} cm^{-3}$. The substrate is heated to provide at least the reaction activation energy. With these films, the photovoltaic cell comprising glass / SnO$_2$:F / TiO$_2$ / CH$_3$NH$_3$PbI$_3$ / CuS / Mo can be prepared as shown in figure 2.1-a. Figure 2.1-b presents an optical microscope image of a multilayer film comprising Na Glass substrate, SnO$_2$:F transparent conducting oxide, TiO$_2$:Ni, and Ni. This sample was annealed at 270°C. The image reveals significant porosity with a length scale of 10s of microns. In films grown at higher temperature, this microstructure was less pronounced.

Photo-generated electrons in the perovskite absorber layer would be injected into the embedded TiO$_2$ wires that would transport these electrons to the transparent-conducting SnO$_2$:F electrode [67]. Photo-generated holes would be injected into the inorganic CuS hole layer with transport to the Mo electrode.

Since the TiO$_2$ lies above the perovskite absorber layer, it must have high transparency in the visible range while maintaining sufficiently high electron conductivity.
Figure 2.1: Figure 2.1-a is a schematic of proposed Perovskite Solar Cell layer structure. Figure 2.1-b is an optical microscope image of SPEED-grown TiO$_2$ thin film.
Borofloat glass substrates were cleaned using a heavy-duty degreaser, acetone, and de-ionized water. Then, the substrate was mounted on a heater set to the growth temperature between 270°C and 400°C with ~10°C uncertainty. The SPEED version used here is vapor phase SPEED (VP-SPEED). The water-based precursor is nebulized into 10 to 20 µm droplets, which are sprayed onto the substrate. Growth rate exceeds 200 nm per minute. Reaction byproducts are decomposed and volatilized. A computer controls the position of the spray nozzle on a 2-axis CNC machining table. Five spray cycles were made for each sample. Annealing was performed for 10 minutes in a tube furnace preheated to 450°C. Scanning electron microscopy (SEM, Zeiss ULTRA-55) with Energy Dispersive Spectroscopy (EDS) was performed at the UCF Materials Characterization Facility (MCF). For cross-sectional SEM imaging, samples were cut and their edges polished. Asymmetric out-of-plane X-ray Diffraction (XRD) was measured using PANalytical Empyrean. The incidence angle with respect to the sample plane was 1 degree. The X-ray wavelength was Cu K-Alpha 1, 0.1540598 nm. Optical microscopy was performed using a Nikon Labophot Microscope.

In-plane resistivity was measured using a 4-point probe. Resistivity normal to the plane was determined by measuring the I-V curve via a 1cm² evaporated Ni contact. A Cary 500i spectrophotometer was used to measure film transmittance spectra using bare substrates in the reference channel.
2.2 Experimental Methods for Smooth TiO$_2$ for Infrared Total Photonics

Our aqueous process is called Streaming Process for Electrodeless Electrochemical Deposition, or SPEED. The substrate is heated to provide at least the reaction activation energy as shown figure 2.2. Figure 2.2 presents a photograph of the SPEED apparatus during deposition. In contrast to chemical bath deposition, SPEED ensures that films grow by heterogeneous reaction only, instead of a mixture heterogeneous and homogenous reactions.

![SPEED Deposition method](image)

Figure 2.2: SPEED Deposition method

For TiO$_2$ growth, precursors comprised 0.1 M titanium chloride in mixed water and organic solvents. The 20%-by-volume deionized water is a solvent, provides surface-attached (OH)$^-$, and is the source of oxygen. Organic solvents ethanol, isopropanol, and methyl propanol serve mainly as complexing agents for Ti ions.
We chose silicate glass and aluminum-coated glass surfaces on which to grow films. Both surfaces are hydrophilic and readily anchor hydroxyl ions that serve as nucleation sites. The \((OH)^-\) ions attract positively charged complexes to initiate electrochemical reaction, without external field or electrodes. The heterogeneous chemical reactions are:

\[
[Sub](OH)^- + [TiL_n]^{p+ (4+)} \rightarrow [Sub]Ti(OH)^{3+} + (L_n)^p \uparrow \tag{2.1}
\]

\[
[Sub]Ti(OH)^{3+} + 3(OH)^- \rightarrow [Sub]TiO_2 + 2H_2O \uparrow \tag{2.2}
\]

where \(L\) is the ligand supplied by the organic solvents, \([Sub]\) is the heated substrate, \(p\) is the charge of the ligand \(L\), and \(n\) is the number of ligands involved in the \(Ti/L\) coordination. Up arrows indicate decomposition and evaporation of reaction byproducts. Substrate temperature must exceed the heterogeneous reaction activation energy and be sufficient to eliminate byproducts. Site regeneration follows

\[
[Sub]TiO_2 + (OH)^- \rightarrow [Sub]TiO_2(OH)^- \tag{2.3}
\]

The freshly attached \((OH)^-\) initiates the next growth cycle. Lateral grain growth competition leads to nanoparticle morphology.

Asymmetric out-of-plane X-ray Diffraction (XRD) was measured using PANalytical Empyrean system. The XRD patterns were measured using CuK\(\alpha\) radiation \(\lambda = 0.154059\) nm at 45 kV and 40 mA with \(2\theta\) scanned from 20 to 65 deg. The incidence angle with respect to the sample plane was 1 degree.
These data were compared with reference spectra for Anatase, Rutile and Brookite from International Centre for Diffraction Data (ICDD) PDF2 Release 2015.

Scanning Electron Microscopy (SEM) was performed using Zeiss ULTRA-55. Cross section samples were prepared by standard polishing technique using Allied Multiprep Polishing System-12 and briefly etched in 10:1 Buffered Oxide Etch solution to enhance image contrast. SEM imaging was done using Zeiss Cross Beam Field Emission Gun (FEG) SEM. Samples for SEM imaging were sputter coated with 2 nm Iridium (Gatan Precision Etching Coating System). Atomic Force Microscopy (AFM) of the TiO\textsubscript{2} films was performed using (DI) Digital Instruments 5000U. Data visualization and analysis used the Gwyddion program. A Cary 500i spectrophotometer was used to measure film-on-glass transmittance spectra. A bare glass substrate was used in the reference channel.

2.3 Experimental Methods for SPEED BaTiO\textsubscript{3} for Pyroelectrics

Self-assembled nano-crystalline BaTiO\textsubscript{3} films on stainless steel foil substrates, were grown by the water based Streaming Process for Electrodeless Electrochemical Deposition (SPEED). The nebulized aqueous precursor impinges on a substrate heated to \(~300^\circ\text{C}\), giving growth rate exceeding 200 nm/min. Substrate heating provides reaction activation energy and decomposes/volatilizes reaction byproducts at temperatures well below those required for spray pyrolysis. Grown films were subsequently annealed in the range 500-600°C for 1-2 hours. Figure 2.3 presents photograph of experimental set up for photoresponse measurements. Physical characterization by Scanning Electron Microscopy (SEM, Zeiss ULTRA-55 FEG) and asymmetric out-of-plane X-ray Diffraction (XRD, PANalytical Empyrean #2) were performed at the UCF Materials Characterization Facility (MCF). The Cu K\textsubscript{α} X-ray beam at 0.1540598 nm was incident at a fixed angle of 1 deg.
Figure 2.3: Photograph of experimental set up for photoresponse measurements. A spring contact connects to the carbon-black coating on the BaTiO$_3$ thin-film. The amplifier is constructed on the back side of the printed-circuit board. The arrow points to the spot of a red alignment laser which indicates the position illuminated by the pulsed CO$_2$ laser.
In preparation for photo-response measurements, a 4µm thick BaTiO$_3$ film on 75µm thick stainless steel foil substrate, which had been annealed 2 hours at 500°C, was coated with a conducting carbon-black IR absorbing layer.

A gold-coated spring contact was made to the carbon black layer and connected to the inverting input of a type 741 operational amplifier, as shown in figure 2.3. The substrate and non-inverting input of the op-amp were grounded. A 1 MΩ feed-back resistor was connected between inverting input and op-amp output, which was connected to a digital oscilloscope.

Photoresponse was excited using a pulsed CO$_2$ laser (MTL-3 Mini TEA, Edinburgh Instruments) at 10 micron wavelength. Pulse duration was approximately 100 ns and repetition rate was 2 Hz. The peak power was not precisely measured, but the power in a single pulse was sufficient to cause a noticeable change in the color of liquid-crystal thermal paper designed for 25-30 deg operation (Edmund Optics R25C5W). The color change was similar to that caused by a 100 mW CW CO$_2$ laser, allowing an estimate of 50 mJ per pulse. This value is the same order as that specified in the laser manual.

The precursor solution included Ba(NO$_2$)$_2$ and TiCl$_2$ as the source of Ba and Ti ions. The likely reaction is

\[
(BaTi) - Ligand^{(p+6)} + substrate - (OH)^- \rightarrow substrate - BaTi(OH)^{5+} + Ligand^{(p)}(gas)
\]  
\[
(\text{substrate}) - BaTi(OH)^{5+} + 5(OH)^- \rightarrow (\text{substrate}) - BaTiO_3 + 3H_2O(gas)
\]  
(2.4)

(2.5)
The superscript $p$ is the ligand charge, which may be positive, negative, or zero. Thirty spray cycles were applied at 300°C, and then the films were annealed at 450 or 500°C for 2 hours. Subsequent annealing at 600°C was also performed on some samples. A purpose of the post growth anneal was to decompose and volatize residual byproducts and recrystallize the formed nanoparticles based film.

Scanning electron microscopy (SEM, Hitachi SU70) with energy dispersive X-ray spectroscopy (EDX) was used for surface and cross sectional imaging and elemental analysis.

Asymmetric out-of-plane X-ray diffraction (XRD, PANalytical Empyrean) was performed to determine crystalline phases. The CuKα X-ray beam at 0.1540598 nm was incident at a fixed angle of $\omega = 10$ deg. The scattered intensity was recorded as a function of scattering angle $2\theta$. For photo-response measurements, a $\sim 1 \mu m$-thick film on stainless steel substrate was coated with a $\sim 30 \mu m$ thick conducting carbon-black IR absorbing layer. Sample area was 1 to 4 cm$^2$. Photo-response measurements were made using two different IR sources and amplifiers.

System A comprised a CO$_2$ laser (Access Lasy 3S) with $\sim 100$ mW CW power, which was mechanically chopped at 70 Hz. Laser spot size on the sample was $\sim 1-2$ mm. A spring contact was made to the carbon black layer and connected to the inverting input of a type 741 operational amplifier with 1 MΩ negative feedback resistor. The substrate and noninverting input of the op-amp were grounded. There was no electronic filtering. A digital oscilloscope averaged 64 traces of the photoresponse.

System B measured pyroelectric current based on the dynamic capacitance method [70]. A broadband black-body source chopped at 1 Hz illuminated the sample with $\sim 13$ mW/cm$^2$. A low-noise high-gain pre-amplifier converted the pyroelectric current into a voltage with gain of 1V/100 pA. The signal was filtered by a 7.2 Hz low-pass. A digital oscilloscope averaged 16 traces.
Ferroelectric hysteresis was measured by the Sawyer Tower method [71]. A sawtooth voltage $V_x$ was applied across the series combination of a capacitor made from our $\sim 1 \ \mu m$ thick film with area $A=1 \ cm^2$ and a sense capacitor $C_s=1 \ \mu F$. Since most of $V_x$ is dropped across the much smaller sample capacitor, the electric field in the film is $V_x$ divided by film thickness. The polarization ($P$) is determined by the extraneous charge on the sample capacitor, which is the same as that on the sense capacitor, so that $P = (C_s/A) \ V_y$, where $V_y$ is the voltage dropped on $C_s$. 
CHAPTER 3: RESULTS AND DISCUSSION OF ROPY TiO$_2$ FOR SOLAR CELLS

Figure 3.1 presents top-view, and cross sectional SEM images of the Ropy TiO$_2$ film. In top view, the films have a ropy texture, with large surface area.

![SEM images for TiO$_2$:Ni film. Figure 3.1-a is a top view. Figure 3.1-b is a cross section.](image-url)
The cross-section reveals the film thickness to be \( \sim 1 \mu \text{m} \). Energy Dispersive Spectroscopy (EDS) confirms the presence of Ti, O, and Ni in the film. There is 16% as much Ni as Ti in the film, which agrees with the mix in the precursor.

Figure 3.2-a presents example XRD data for one of the films. The growth temperature was 300°C and sample was not annealed. XRD spectra were collected for all samples, and their XRD patterns were found to be independent of growth temperature and annealing. The spectra were compared with reference spectra for Anatase, Rutile and Brookite (shown) from International Centre for Diffraction Data (ICDD) PDF2 Release 2014. The phase was identified as Brookite (orthorhombic).

Cross sections of all grown samples reveal the layer thickness dependence on growth conditions. Figure (3.2-b, top) presents a plot of film thickness as a function of growth temperature. Post-growth annealing caused no significant difference in the film thickness, so results for annealed and unannealed samples are averaged. For the same 5 spray cycles, the thickness decreases by a factor of \( \sim 2 \) on increasing growth temperature from 270°C to 400°C. Figure (3.2-b, bottom) presents the measured in-plane resistivity and sheet resistance as a function of growth temperature. The two curves are not exactly proportional because of differences in film thicknesses for different samples. The resistivity increases with growth temperature and covers the range \( 40 \text{ to } 200 \, k\Omega \cdot \text{cm} \). In plane conduction less important for solar cell function than vertical transport, since photo-electrons must pass vertically through the TiO\(_2\) layer to reach the Fluoride doped Tin Oxide (FTO) transparent conducting top contact. If the resistivity to vertical transport is the same as the horizontal, we would expect a resistance for \( 1 \, cm^2 \) cross section and 1 micron thickness to be in the range \( 4 \text{ to } 20 \Omega \). However, the vertical resistance through a Ni pad of this area was found to be 4500\( \Omega \), presumably due to high contact resistance.
Figure 3.2: Figure 3.2-a is a typical XRD spectrum of SPEED-grown with Brookite phase reference data. Figure 3.2-b is a resistivity, sheet resistance, film thickness, and optical transmittance at 500 nm wavelength versus growth temperature.
The optical transmittance spectra of the films all have the same rapid drop at wavelengths below 388 nm, which corresponds to fundamental absorption above the 3.2 eV TiO₂ bandgap. Transmittance increases with growth temperature as shown in figure (3.2-b, top), which may be reasonably attributed to a combination of the correlated thickness decrease and $R_S$ increase.
CHAPTER 4: RESULTS AND DISCUSSION OF SMOOTH TiO$_2$ FOR INFRARED TOTAL PHOTONICS

We designate the film considered here as Smooth SPEED TiO$_2$ to distinguish it from a highly structured film grown by a different SPEED recipe studied earlier[10], which was designated Ropy SPEED according to its observed morphology. Figure 4.1 presents XRD data and shows that Smooth SPEED TiO$_2$ on Al-coated glass is Anatase phase. (Growth of TiO$_2$ on a conducting substrate was motivated by certain Long-Wave Infrared (LWIR) applications of interest [50]). This XRD spectrum also shows definitively identified aluminum peaks from the metal coating on the substrate.
Figure 4.1: XRD Smooth TiO$_2$ film and reference spectra

Figure 4.2 presents SEM images of the top surface of the film.
The surface is mostly smooth and featureless, except for small scattered blisters, mostly with lateral dimensions of order $\sim 1\mu m$. These films are suitable for long-wave infrared (LWIR) applications, because those wavelengths ($\sim 10\mu m$) are much larger than the surface defects.

Figure 4.2: SEM image of Smooth SPEED TiO$_2$ film surface
Figure 4.3 presents SEM cross sections of the film. Substrate, aluminum, TiO$_2$ layers, and 200 nm scale bar are indicated. The thicknesses of the Smooth SPEED TiO$_2$ film is approximately 130 nm.

![SEM cross-sectional image of Smooth SPEED TiO$_2$ film.](image)

Figure 4.3: SEM cross-sectional image of Smooth SPEED TiO$_2$ film.
Figure 4.4 presents an AFM image of the Smooth SPEED film sample. The surface is characterized by sparse, isolated, and narrow peaks of \(~200\ \text{nm}\) amplitude and \(~1\ \mu\text{m}\) wide base on a smooth background of \(~10\text{nm}\) roughness, which is consistent with the SEM image 4.2.

Figure 4.4: AFM images of Smooth SPEED TiO$_2$. 
Figure 4.5 presents a UV-visible transmission spectrum for a portion of the film where the substrate was uncoated with Al. The values for the refractive index of anatase at 2.5 eV is \( \sim 3 \) for unpolarized light and random crystallite orientation [72].

Figure 4.5: Transmission spectrum of Smooth SPEED TiO\(_2\).
Thus, the Fresnel reflectance $R = (n - 1)^2/(n + 1)^2$ would be 25% at 500 nm wavelength. The peak transmittance at 3 eV is 80%, i.e. larger than expected. This may be due to optical interference within the high index thin film. At longer wavelengths, where the index is smaller, scattering weaker, and fundamental absorption out of range, we might expect that the transmittance should increase, but the opposite is observed. The optical thickness $n \times d$ of the film is $\sim 390$ nm, which exceeds half the 500 nm wavelength, so that interference in the film can be important. This is supported by the observation of one full Fabry-Perot oscillation period in our transmittance spectrum of a 210 nm thick evaporated TiO$_2$ film in the same spectral range.

For the thinner Smooth-SPEED film, Fabry-Perot fringes should be separated by 1.5 eV. A weak bump appears in figure 4.5 at 1.2 eV below the 3 eV peak. This might be the Fabry-Perot fringe, since the high frequency side of the 3 eV peak is truncated by fundamental absorption.

Figure 4.6 presents analysis of the UV-visible transmittance spectrum to determine the band gap $E_g$ of the film. The absorption coefficient $\alpha$ was calculated according to Beers law $\left(1/d \ln[1/T]\right)$ from film thickness $d$ and the transmittance spectrum $T$, which was first normalized to remove the Fresnel reflectance factor. Anatase is an indirect gap semiconductor [72], whose absorption edge depends on absorption and emission of phonons, according to [73].
Figure 4.6: Determination of band gap of Smooth SPEED TiO$_2$ film.

\[
\alpha_e + \alpha_a = C_e(T)(h\nu - E_g - E_p)^2 + C_a(T)(h\nu - E_g + E_p)^2,
\]  

(4.1)
where the coefficients depend on temperature-dependent phonon statistics. A plot of the square root of equation 4.1 gives a curve with two intersecting straight-line segments of different slope. The extension and intersection with the zero-absorption axis define the energies \( E_g \pm E_p \). The determined values are given in figure 4.6. Often it is the experimental values of \( \alpha h \nu \) that are equated to equation 4.1 [74]. A plot by this method analysis looks very similar to figure 4.6, but it gives the somewhat larger values \( E_g = 3.21 \) eV and \( E_p = 0.21 \) eV. The latter value is more within the range 3.2-3.35 eV reported for the indirect gap of Anatase [72, 75, 76, 77, 78, 79]. The phonon energy seems rather high, because the onset of optical phonon absorption appears below 0.1 eV, as shown figure 4.7. At high energies, the absorption coefficient increases less rapidly with photon energy.

Its functional dependence is expected to be \( A(h \nu - E'_g)^{1/2} \), where \( E'_g \) is the direct gap. We find \( E'_g \) to be about 4.2 eV by extrapolating a plot of to zero absorption. This value is higher than the expected value 3.8 eV [72]. However the portion of the spectrum that follows a square root function is very narrow, and it falls in a region of very low transmittance, so that the value of \( E'_g \) thus obtained has high uncertainty.

Figure 4.7 presents infrared spectra of the real and imaginary parts for the complex refractive index \( n + i \kappa \), where \( n \) is the refractive index and \( \kappa \) is the extinction coefficient. These spectra compare well with previously published results for films deposited by Atomic Layer Deposition (ALD) [44] and sputtering[45] . For comparison we present our data for an evaporated film of the more common dielectric used in planer photonic devices SiO\(_2\). In contrast with SiO\(_2\), which has a strong and sharp extinction peak near 9 \( \mu \)m wavelength, the extinction of TiO\(_2\) is relatively weak out to at least 11 \( \mu \)m wavelength. Thus, TiO\(_2\) would be more attractive than SiO\(_2\) as an index contrast material for Long-Wave Infrared LWIR photonic planer waveguides due to its comparatively low loss. Moreover, in contrast to the derivative-like feature in the refractive index of SiO\(_2\), the index of TiO\(_2\) changes monotonically.
This is particularly important in resonance absorption structures, such as metal-insulator-metal plasmonic absorbers, where derivative-like dispersion spectra give multiple absorption peaks that complicate the spectrum and interfere with spectral selectivity[48, 80].

Figure 4.7: Infrared complex index spectra for Smooth SPEED TiO$_2$ and evaporated SiO$_2$ films. The shaded region indicates LWIR wavelengths.
The onset of extinction beyond 12 $\mu$m wavelength suggests optical phonon energies less than 0.1 eV. This is agreement with optical phonon frequencies determined from far-infrared reflectivity [81]. Thus, the phonon energy determined in from figure 4.6 is surprisingly large, unless fundamental absorption at the indirect gap is a two phonon process.
CHAPTER 5: RESULTS AND DISCUSSION OF EVAPORATED TiO$_2$

The XRD patterns were measured at a grazing incidence angle of 1° using CuK$_\alpha$ radiation $\lambda=1.54059$Å at 45 KV and 40 mA with a pattern recorded from 20° to 65° of 2$\theta$. Evaporated TiO$_2$ is a Anatase phase, as shown in figure 5.1, after annealing. Before annealing, XRD had no sharp peaks, so as evaporated film was amorphous.

![Evaporated Annealed TiO2](image)

![Anatase](image)

Figure 5.1: XRD measurement of Evaporated TiO$_2$ on Glass.
UV-visible transmittance measurements are analyzed to determine the band gap of each film. The absorption coefficient is related to the direct band gap and the incident photon energy by $\alpha h\nu = C(h\nu - E_g)^{\frac{1}{2}}$, Absorption coefficient was calculated according to Beers law, $\alpha = \left( \frac{1}{d} \right) \ln \frac{1}{T}$, Where $d$ is film thickness. The band gap of evaporated TiO$_2$ is found to be about 4.0 eV, which appears somewhat high, though not as much so as for Smooth TiO$_2$ as shown in figure 5.2.

Figure 5.2: Band gap of E-beam evaporation.
The largest transmission of E-beam evaporation TiO$_2$ at 380 nm is 93% as figure 5.3. This film shows a complete Fabry-Perot oscillation. The separation in wave numbers $\Delta\tilde{\nu}$ is about 6700 cm$^{-1}$. The separation should equal $\frac{1}{(2nd)}$, which for $n=3$ is 7200 cm$^{-1}$. The observed value is smaller because fundamental absorption crops the short wave peak.

Figure 5.3: Transmission of E-beam evaporation.
Evaporated TiO$_2$ is rough with large surface area as shown in figure 5.4.

Figure 5.4: SEM Top View of Evaporated TiO$_2$. 
SEM cross-section of Evaporated TiO$_2$ has thickness $\sim$ 230 nm as shown in figure 5.5.

![SEM cross-sectional image of Evaporated TiO$_2$ film.](image)

**Figure 5.5:** SEM cross-sectional image of Evaporated TiO$_2$ film.

Atomic Force microscopy (AFM) images of the TiO$_2$ films obtained using (DI) Digital Instruments 5000U, which operated in the constant mode by putting a sample under Silicon with Aluminum reflex coating tip, scanning in Air. The spring constant of the Cantilever was 40 N/m and resonant frequency was 300 kHz. Actual resonant frequency was 280.69 kHz and actual spring constant was 31.01 N/m. Scan size was 30μ m. AFM analyzed of all samples (Ropy, Smooth, and Evaporated). By using (DI) Digital Instruments 5000U. Data visualization and analysis used to Gwyddion program.

An AFM image for the evaporated sample is presented in figure 5.6. It shows the surface comprises flat terraces with a small number of defects. Scan area of AFM image 30 μm x-axes, 15 μm y-axes. The vertical dimension range is 0.13 μm.
Figure 5.6: AFM 3-Dimensions of Evaporated TiO$_2$ Sample.
CHAPTER 6: RESULTS AND DISCUSSION OF SPEED BaTiO$_3$ FOR PYROELECTRICS

Figure 6.1-a presents a top-view SEM image of a BaTiO$_3$ film grown by SPEED. The film comprises a sea of nano crystals, in which are embedded micron-sized single-crystal islands. Figure 6.1-b presents a cross sectional SEM image of the film, revealing a 2 $\mu$m film thickness.
Figure 6.1: SEM images of SPEED-grown BaTiO$_3$ films. Figure 6.1-a is a top view revealing a mix of nano- and micro-meter scale crystals. Figure 6.1-b is a SEM cross section.

Figure 6.2 presents XRD spectra for two different annealing temperatures indicated.
Reference spectra are also presented, showing that the sample annealed at 500°C, 2 hours is hexagonal, while the sample annealed at 600°C, 1 hour has converted to the non-pyroelectric cubic phase. Notable differences are that the latter is missing the peaks evident in the spectrum of the former near 18, 37, 49, 59, and 69 degrees. Both measured spectra contain peaks that are absent from either reference spectrum. Extra peaks at 24, 34 and 42.5 deg are stronger in the 600°C annealed sample, which also contains an extra line at 46 deg absent for the 500°C sample. These extra peaks suggest the presence of one of the other BaTiO$_3$ phases that tends to increase on annealing at 600°C.
Figure 6.3-a presents an oscilloscope screen shot with a response transient for a sample annealed at 500°C for 2 hours. The decay time is about 5 ms. There is strong electrical disturbance due to the laser pulsing circuit at times earlier than shown, with the result that when the laser beam is shuttered, there is a negative signal of about the same magnitude in the region indicated by the boxcar gate (channel 2).
Figure 6.3: Figure 6.3-a is an oscilloscope screen shot showing the transient pyroelectric response to the laser pulse on channel 1. Channel 2 shows the gate of a boxcar average. Figure 6.3-b is a peak signal as function of position on the absorbing area of the BaTiO$_3$ device.
Thus, the peak photo response is about 400 mV, or 8 V/J for the sample and amplifier combination. This signal was 10x stronger than for a similar sample annealed half as long (1 hr at 500°C). Figure 6.3-b presents a measurement of photo-response as a function of position on the 1 cm² area of the carbon-black coating.

The sample was translated using an x-y stage. The point (0, 0) corresponds to the center of the sample. The x-scan was performed at y = zero and vice versa. The photo response appears to be fairly uniform across the active area, dropping steeply the edges of the carbon-black square.

Note that the signal strength in figure 6.3-a is smaller than in figure 6.3-b which was measured about one week earlier. The photo-response was found to decay over time, presumably due to the absorption of water vapor from the air.
Figure 6.4-a presents a plot of 4 measurement points collected over a period of 3 weeks. The decay has a characteristic time of 1 week. Baking at 100°C restores the response, as shown in figure 6.4-b.

Figure 6.4: Figure 6.4-a is decay in peak photo-response with time for sample kept in air. Figure 6.4-b is a recovery of signal on baking at 100°C.
Poling experiments were performed on the sample on day 13 as shown in figure 6.4, when the signal had decayed almost to zero. This was achieved by applying a DC voltage between top and substrate surfaces of the sample with the op-amp power turned off. Voltages of 0.5 and 1 V applied for 2 minutes produced no effect. Voltages of 2 V applied for 2 minutes produced a negative signal with $-40$ mV peak amplitude, and this signal had the opposite sign from the original signal of the unpoled sample. The signal achieved by poling decayed in $\sim 1$ min. The sample could be poled again, and the behavior was repeatable. When poling of the opposite sign was applied, the photoresponse was positive.
Figure 6.5 presents a high resolution top-view SEM image of another BaTiO$_3$ film. We refer to this film as BaTiO$_3$ even though subsequent analysis shows that it is only partially composed of this compound, in contrast to our previously reported films [69], in which BaTiO$_3$ was the dominant crystalline phase. The SEM image reveals grains with a size distribution of about 50 to 200 nm. Additional SEM images reveal a fine grained matrix that contains cracks.

Figure 6.5: High resolution SEM image of SPEED-grown BaTiO$_3$ thin film.
Figure 6.6 presents an SEM image of the surface at the edge of the carbon-black coating at low magnification. The BaTiO$_3$ film has a white crust, which EDX shows to be up to 35% Cl. This indicates that the Cl byproduct was incompletely volatized, and it suggests the need for longer and hotter anneals. EDX on the carbon-black film shows almost exclusively carbon. Cross sectional SEM imaging shows the BaTiO$_3$ and C-black films to be $\sim 1$ and $\sim 30 \mu$m thick, respectively.

Figure 6.6: SEM image of surface of SPEED-grown BaTiO$_3$ film with carbon-black coating (lower region).
Figure 6.7 presents powder XRD reference data for the various phases of BaTiO$_3$.

The signal-to-noise ratio is sufficiently low in the XRD data for our thin films that the non-hexagonal phases would be difficult to distinguish from each other. The XRD pattern for the hexagonal phase is much different than the others.
The relative intensities of XRD peaks for nano-crystalline thin films can differ significantly from those of the powder due to interactions with the substrate.

Figure 6.8 compares our measured XRD data for films annealed after growth at the three different temperatures 450, 500, and 600°C.

Figure 6.8: XRD spectra of SPEED-grown BaTiO$_3$ thin films. Annealing temperatures in degrees C and the reference positions of XRD peaks for cubic and hexagonal phases are indicated.
The annealing times were one, 2, and 4 hours respectively. At a given annealing temperature, the XRD spectrum was independent of annealing time in the range 1-4 hours. The substrate spectrum was taken on the backside of the substrate and reveals substrate lines at \( \sim 44 \) and 51 degrees. These substrate lines appear as artifacts in the XRD of the two lower-temperature annealed samples. The expected peak positions for hexagonal and cubic BaTiO\(_3\) are indicated. In all three spectra, the dominant peaks do not belong to any of the 5 BaTiO\(_3\) phases.

In the two lower temperature annealed samples there are several weak peaks that could correspond to the hexagonal phase, but they could with equal probability be assigned to phases of residue TiCl\(_2\), TiCl\(_3\), or TiCl\(_4\), and the strongest peaks in these spectra belong with high confidence to a hexagonal phase of TiCl\(_3\), which itself could be pyroelectric. In the 600\( ^\circ \)C annealed sample there are a few small XRD peaks that can be identified with one of the non-hexagonal phases. Thus it appears that the films are a mixture of BaTiO\(_3\) nano-crystals in a matrix of some other polycrystalline material.

The transition from hexagonal to one of the other phases with increasing annealing temperature agrees with our previous observation on a different set of BaTiO\(_3\) films [69], which were more nearly pure BaTiO\(_3\). Scherrers formula [82] indicates that the lower bound on the mean size of the ordered crystalline domains of BaTiO\(_3\) is \( \sim 20 \) nm.
Figure 6.9-a presents photoresponse data for a carbon-black coated sample annealed at 500° C for 2 hours using measurement system A.

(a)

Figure 6.9: Figure 6.9-a is an oscilloscope screen shot showing pyroelectric response to 100 mW CO$_2$ laser chopped at 70 Hz, amplifier A. Figure 6.9-b is a response to 13 mW/cm$^2$ broad-band IR source chopped at 1 Hz, amplifier B.

(b)

Figure 6.9: Figure 6.9-a is an oscilloscope screen shot showing pyroelectric response to 100 mW CO$_2$ laser chopped at 70 Hz, amplifier A. Figure 6.9-b is a response to 13 mW/cm$^2$ broad-band IR source chopped at 1 Hz, amplifier B.
The signal comprises a peak when the chopped laser first illuminates the sample, and a smaller and slower signal when the laser is chopped off. The response time is about 1 ms. The responsivity is about 100 mV/W, compared with the $\sim 1 \mu$V/W we previously reported for a more-nearly pure hexagonal BaTiO$_3$ film using a 100 ns pulsed CO$_2$ TEA laser with a similar amplifier [69]. That earlier peak response was reduced by the ratio of laser pulse duration to system response time, $100\text{ns}/5\text{ms} = 20 \times 10^{-6}$. Scaling the earlier result by that ratio gives $\sim 50$ mV/W as the probable response that would be measured for that earlier film using measurement system A. Thus, the present measurement is in agreement with the earlier one.

Figure 6.9-b presents photo-response using measurement system B. The signal is now symmetric with respect to the on-off chopping of the source. The peak output voltage was 0.5 V. The response has a 400 ms time constant, which is slower than expected for the 7.2 Hz filter cut-off frequency. With system B, we achieve responsivity of 40 V/W.

Previously we reported the appearance of a short-lived (1 min) photoresponse that appeared after poling the film with a few volts bias at room temperature. The sign of the photoresponse depended on the sign of the applied bias. We now interpret this effect as poling-induced charging of traps, whose spatial distribution depends on the sign of the applied bias. When the sample is hit by infrared, these traps are thermally discharged, giving a current that is comparable to pyroelectric current. As soon as the traps are all emptied, the effect disappears. This effect is known to be especially significant in thin films. [83]
Figure 6.10 presents ferroelectric hysteresis measured on the film studied in ref.[69], which is 5x smaller than the value for bulk BaTiO$_3$ [84]. That film was primarily hexagonal phase, with significant admixture of some other material, which can explain the relative smallness of its polarization compared to bulk tetragonal BaTiO$_3$.

![Ferroelectric hysteresis curve of BaTiO$_3$ thin film.](image)

Refs.[49], and [66] described a design for spectrometer on a chip. The design comprised a broadband IR source, rectangular photonic waveguide, interaction region, dispersion region, and detectors.
In that work, only the waveguide section formed by deep reactive ion etching of silicon had actually been fabricated. The proposed sources were IR LEDs, but integrating these onto the chip had not been considered. An adiabatic taper transduced the photons in the waveguide to surface plasmons polaritons that existed in a gap between the outer side walls of the narrowed waveguide and vertical side walls of border material coated with suitable conductor. We have since demonstrated perfectly conformal conducting SnO\textsubscript{2}:F that can be deposited by SPEED on such sidewalls to support surface plasmons in this interaction region [68]. With Ref.[69] and this paper we have presented first results on a pyroelectric BaTiO\textsubscript{3} detector material, which can be deposited by SPEED, and which in principle can be positioned at the end of a waveguide by standard photolithographic patterning. Clearly, considerable work remains to optimize the film, to develop processes and to design suitable amplifiers.
CHAPTER 7: CONCLUSION AND FUTURE WORKS

We reported physical, electrical, and optical properties of nano-crystalline orthorhombic Ropy TiO$_2$ thin films of micron thickness grown by an aqueous spray deposition method. The proposed Perovskite Solar Cell presented is a type of considerable current interest, which can benefit from our TiO$_2$ layer as electron collector, and which can be entirely fabricated by our aqueous process. Work to complete and test such a cell is ongoing.

We presented optical and morphological properties of Smooth TiO$_2$ film grown by aqueous spray deposition. The crystalline phase is Anatase. Defects that are present are much smaller than long-wave infrared wavelengths (LWIR) and should not contribute significantly to electrodynamic losses. Low dispersion in the long-wave infrared is attractive for plasmonic resonant absorbers in applications that require spectral selectivity.

BaTiO$_3$ thin films were grown by an aqueous spray deposition process at moderate substrate temperatures. Two phases were observed depending on annealing temperature. The hexagonal phase demonstrated a pyroelectric infrared photoresponse. An aqueous spray deposition process grew BaTiO$_3$ thin films. Two phases were observed, depending on annealing temperature, in a matrix of some other material.
LIST OF REFERENCES


62


