A Parametric Study of Economical Energy Usage in Freeze Tunnels

Summer 1980

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A PARAMETRIC STUDY
OF ECONOMICAL ENERGY
USAGE IN FREEZE TUNNELS

BY

MARC A. HARRISON
B.S., United States Naval Academy, 1973

THESIS

Submitted in partial fulfillment of the requirements
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ABSTRACT

An investigation into economical energy usage in freeze tunnels was conducted. Freeze tunnels are commonly used in the food processing industry to freeze products, and in some cases may use large amounts of electricity. An actual freeze tunnel was observed and modeled on a computer.

A parameter study was conducted. The results of the parameter study indicate the efficiency and energy costs in freeze tunnels may vary widely. Important parameters included the Nusselt number, air temperature, and the ratio of fan work divided by the useful refrigeration effect. Although no single set of optimum conditions were found, methods for improving the effectiveness of freeze tunnels, both in existing and future designs, were discussed. It was also concluded that the ratio of fan work to the freeze tunnel's useful refrigeration effect was a dominant factor in the energy cost of operating a freeze tunnel.
ACKNOWLEDGEMENTS

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CHAPTER 1

INTRODUCTION
1.1 INTRODUCTION

An investigation into economical energy usage in freeze tunnels was conducted. The primary objective of this investigation, was to determine the effects of the important design and operating parameters on energy consumption in these devices. Freeze tunnels are commonly used in the food processing industry to rapidly freeze or reduce the temperature of food products. Rapid cooling is often required to preserve food quality and to meet production goals. In some cases, the cost of operation of freeze tunnels is a small part of the cost of the entire food processing operation [1]. But, as energy costs continue to rise efficient energy usage will become more important. In other cases energy consumption in freeze tunnels is already a large part of the energy consumed in the entire operation. One study estimated for a medium sized citrus juice concentrate processing plant, about 25% of the total energy costs, of about $1.4 \times 10^6$ dollars per season, was due to freeze tunnel electricity consumption [2].

The parameter study was accomplished with a computer model of a freeze tunnel. The computer model was based on an actual freeze tunnel that was available for observation. Measurements of the actual freeze tunnel's typical operating conditions were made and compared with predictions of the computer model. The computer model was initially pro-
grammed to simulate actual tunnel operating conditions as closely as possible. After the validity of the model was demonstrated, important parameters were varied from the actual conditions measured for the observed tunnel. The effects of the parameter variations on the freeze tunnel's effectiveness was then evaluated.

1.2 FREEZE TUNNEL DESCRIPTION

Freeze tunnel designs may vary with usage, capacity, food product, and manufacturer. The tunnels studied in this report are used to rapidly reduce the temperature of orange and grapefruit juice concentrate just after it is canned. Parameters that affect energy consumption in the tunnel observed are assumed to have similar effects in freeze tunnels in general. In any freeze tunnel, energy is consumed primarily by the fans and the refrigeration units. Figure 1 is a simple sketch of the freeze tunnel observed with approximate dimensions.

The freeze tunnel is used by a citrus concentrate plant in Central Florida. It is located inside a large building which shields it from environmental extremes. Right circular cylindrical cans of citrus concentrate enter the freeze tunnel on a mesh conveyor belt. In general, the cans stand upright and are packed tightly together. Refrigerated air is blown between the cans by large fans to maintain a high rate of heat transfer and short freezing times. Although, the conveyor belt is driven by a single speed motor and reduction gear, the conveyor belt may be stopped for short periods of time as required by events in other parts of the production line. Ideally, the cans exit the tunnel, on the conveyor belt, simultaneously with the desired freezing time. The
Fig. 1. Freeze Tunnel Sketch

SIDE VIEW

150 ft

END VIEW

15 ft

20 ft

Conveyor Belt

Evaporator

Cutout

Fan

Interior Walls

Conveyor Belt
tunnel's evaporators are supplied with ammonia refrigerant by a large two-stage, vapor compression plant. There are two large doors on opposite sides of the freeze tunnel. The ceiling and walls are insulated by 6 inches of polyurethane insulation, encased in metal. The floor is a cement slab. The quantity and type of any insulation in the floor could not be determined. The tunnel contained 10 fans rated at 10 horsepower each and 8 evaporators.
CHAPTER 2

THEORY
2.1 THE FREEZE TUNNEL COEFFICIENT OF PERFORMANCE

In this report, the coefficient of performance (COP) for the freeze tunnel is used to estimate the effectiveness of the tunnel. The coefficient of performance is defined [3] as

\[ \text{COP} = \frac{\text{refrigeration effect}}{\text{net work input}} \]

In this case the useful refrigeration effect is the rate of heat removal from the concentrate, \( q_c \). The net work input is the sum of the work of the fans, \( W_f \), and the work of the compressors, \( W_c \). Thus, the coefficient of performance for the tunnel \( \text{COP}_t \), becomes

\[ \text{COP}_t = \frac{q_c}{W_c + W_f} \]  

(1)

The value of \( q_c \) is obtained by calculating the rate of change of concentrate enthalpy in the tunnel and is discussed further in the next two sections. The value of \( W_f \) is estimated from the fan ratings. The value of \( W_c \), in this case, must be determined indirectly because the refrigeration plant supplies several loads besides the freeze tunnel. Therefore, \( W_c \) is estimated by using an energy balance to calculate the refrigeration load.

An energy balance is performed as follows. Steady state operation is assumed so the time rate of change of the stored energy equals zero. The concentrate packing material is ignored. The energy balance then becomes

\[ q_1 = q_c + q_{\text{trans}} + q_{\text{inf}} + q_f \]  

(2)

where
\[ q_L = \text{total refrigeration load} \]
\[ q_c = \text{net rate of energy removal from concentrate} \]
\[ q_{\text{trans}} = \text{transmission heat gain due to conduction and convection to the environment} \]
\[ q_f = \text{rate of energy addition due to fans} \]

Transmission heat gains are calculated with the following equation [4].

\[ q_{\text{trans}} = UA (T_o - T_i) \]

where

- \( U \) = air to air heat transfer coefficient
- \( A \) = area of exposed surface
- \( T_o \) = outside air temperature
- \( T_i \) = average air temperature in refrigerated space

The value of \( U \) for the roof and walls is based on the construction [4]. The value of \( U \) for the cement slab floor is assumed to be 0.1 Btu/hr ft\(^2\)F [5]. A ground temperature of 60°F is assumed. The value of \( T_o \) is based on the summer design dry bulb temperature for Central Florida for the roof and walls. The ambient temperatures were chosen as worst case values to be conservative. Values for \( T_i \) and area were either measured or chosen to correspond to expected operating conditions.

Infiltration heat gains are calculated from [4].

\[ q_{\text{inf}} = 4.5 \text{ (cfm) } \Delta h \]

where

- \( \text{cfm} \) = cubic feet per minute of air infiltrating the tunnel
- \( \Delta h \) = difference in enthalpy between the outside and inside air.
The change in enthalpy was calculated for design summer conditions. The cfm was calculated in two parts: the first part was the cfm due to door openings, and the second part was the cfm infiltrating with the conveyor belt. The cfm due to door openings was calculated using the procedures in ASHRAE [4]. Air is assumed to infiltrate at an average velocity of 75 ft/min. The average cfm is then calculated from the size of the door opening and the fraction of each hour the door is actually open. The second part, air that infiltrates with the food product, is calculated by assuming all air between the cans on the conveyor belt, in the void space, is removed with the cans and replaced by outside air. The volume of the void space and its volumetric flow rate can be measured or specified by operating conditions.

Finally, since the fans are entirely enclosed in the freeze tunnel, their heat addition, in BTU per hour, is given by [4].

\[ q = 2995 \, Hp \]

where

\( Hp = \text{motor horsepower} \)

Once the refrigeration load, \( q_L \), is determined, the required compressor work can be determined from the COP of the refrigerating plant.

\[ W_C = \frac{q_L}{\text{COP}} \]  

Combining equations (1) and (3) results in

\[ \text{COP}_t = \frac{q_C}{(q_L/\text{COP} + W_C)} \]
2.2 THE HEAT REMOVAL RATE FROM CYLINDRICAL CANS OF CITRUS CONCENTRATE

In the case of a freeze tunnel, the useful refrigeration effect is the heat removal rate from the citrus concentrate in the freeze tunnel control volume, \( q_c \), when the packing material is ignored. Calculation of \( q_c \) is complicated by the freezing process of citrus concentrate, the convective boundary condition of the can surfaces, and the substantial temperature gradients that exist in the cans and the tunnel as a result of the rapid freezing process. The best available thermal property data for citrus products has been recently compiled by Chen [6], and this data is currently being evaluated and improved by the Florida Department of Citrus.

A detailed knowledge of the temperature distribution in each can of concentrate versus time is required to mathematically model a freeze tunnel. Knowledge of the temperature distribution is necessary to determine the heat content of each can. Also, the surface temperature of each can is necessary to determine the rate of convective heat transfer from each can to the freeze tunnel environment.

Methods exist to predict temperature distribution changes in freezing problems in general [7,8,9]. Common methods involve assuming a boundary exists between regions of frozen and unfrozen liquids. Each region has appropriate thermal properties and the latent heat is assumed to be evolved at the boundary as it moves through the freezing material. However, as pointed out by Keller and Ballard [9], the freezing process in fruit juice is different.
They considered fruit juice solutions to have the freezing properties of a typical two phase system of ice and solution. In equilibrium, at a given temperature below the freezing point, a given amount of ice exists with a given amount of solution at a certain concentration. Any change in equilibrium temperature alters the amount of ice and solution with a corresponding change in the solution concentration. As the amount of ice and solution changes with temperature, the thermal properties change. Also, the latent heat of fusion for the ice is released or generated over a range of temperatures.

Keller and Ballard calculated values of effective thermal properties over a range of temperatures and citrus juice concentrations. The effective thermal properties, specifically the effective specific heat capacity, \( c_{\text{ef}} \), effective thermal conductivity, \( k_{\text{ef}} \), and density, \( \rho \), include the effects of the latent heat of fusion and any thermal property changes with temperature [9].

Effective thermal property data are used in this investigation. The data chosen correspond to a citrus juice concentration at Brix° 44.8 which is currently a legal standard for Florida orange juice concentrate. Unfortunately, effective thermal property data of concentrate are only available down to temperatures of -20°F and the freeze tunnels considered have been observed producing air temperatures down to about -40°F. Therefore, it was assumed the thermal property data were constant between -20°F and -40°F. The properties are relatively constant with temperature near -20°F. Also temperatures below -20°F were very rarely predicted by the computer and never observed. A summary of the actual data used is listed in
Once the effective specific heat capacity, $c_{ef}$, is known, the heat removal rate from the concentrate can be estimated by integrating

$$q_c = \dot{m} \int c_{ef} \, dT$$

(5)

where

- $\dot{m}$ = concentrate mass flow rate through the freeze tunnel
- $T$ = concentrate temperature

The integrations were accomplished graphically between the average concentrate temperatures at the tunnel entrance and exit. Of course, this method requires established values of both average entrance and exit temperatures.

Another common method, that can be used to calculate $q_c$ is to use Newton's law of cooling,

$$q_c = hA \left( T_s - T_a \right)$$

(6)

where

- $h$ = convective heat transfer coefficient
- $A$ = exposed surface area
- $T_s$ = surface temperature
- $T_a$ = air temperature

This is discussed further in Section 2.3.
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2.3 CONVECTIVE HEAT TRANSFER COEFFICIENT

It is necessary to evaluate the convective heat transfer coefficient, \( h \), to determine the rate of heat transfer from the citrus juice concentrate as a function of time and position in the freeze tunnel.

Whitaker [10] presented a method to calculate \( h \) for flow in packed beds. The packed bed analogy seems appropriate based on observations of the operating tunnel. Although most cans stood upright and were packed tightly, empty gaps and a few cans on their sides were scattered between regions of tightly packed cans.

The method described by Whitaker [10] is briefly presented here. The convective heat transfer coefficient, \( h \), is defined by

\[
q = h \, a_v \, V \, \Delta T_{1n}
\]

where

\[
q = \text{total rate of heat transfer from the packing}
\]

\[
a_v = \text{packing surface area per unit volume}
\]

\[
V = \text{total volume of the packed bed}
\]

\[
\Delta T_{1n} = \text{1on mean temperature difference}
\]

The surface area per unit volume, \( a_v \), is related to the void fraction of the bed, \( \varepsilon \), which is defined as

\[
\varepsilon = \frac{\text{void volume in the bed}}{\text{total volume of the bed}}
\]

The equation is

\[
a_v = (A_p/V_p)(1-\varepsilon)
\]

where

\[
A_p = \text{particle area}
\]
V_p = particle volume

Whitaker [10] showed that the hydraulic radius of the packed bed, R_h, is given by

\[ R_h = \frac{\varepsilon}{a_v} \]

However, the characteristic length of the packed bed, L*, was defined as

\[ L* = 6.0 \, R_h \]

The characteristic velocity, u*, or the average air velocity in the bed, is defined by

\[ u* = \frac{1}{A_{\text{void}}} \int u \, dA_{\text{void}} \]  

(9)

where

- \( A_{\text{void}} \) = cross-sectional void area
- \( u \) = local air velocity

If the bed is uniform, then

\[ u* = \frac{Q}{(\varepsilon A)} \]  

(10)

where

- \( Q \) = air volumetric flow rate through the packed bed
- \( A \) = cross-sectional area of bed

The Reynolds number, Re, the Nusselt number, Nu, and the convective heat transfer coefficient, h, are given by

\[ \text{Re} = \frac{u* L*}{v} \]

\[ \text{Nu} = (0.5 \, \text{Re}^{1/2} + 0.2 \, \text{Re}^{2/3}) \, \text{Pr}^{1/3} \]

\[ h = \frac{\text{Nu} \, k}{L*} \]
The temperature distribution in each can of concentrate must be determined to calculate the can's average temperature, heat content, and surface temperature. The temperature distribution, as a function of time and position in the tunnel, was numerically calculated using an IBM 360.

In this case, the applicable energy equation for heat flow in a cylinder with a convective boundary condition is [11].

\[
\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} + \frac{1}{\alpha} \frac{\partial T}{\partial t} = 0
\]

where

- \( T \) = temperature
- \( \alpha \) = thermal diffusivity
- \( t \) = time

with boundary conditions such that

1) \( T \bigg|_{t=0} = T_i \)
2) \( k \nabla T \bigg|_{\text{surface}} = h(T_s - T_\infty) \)

where

- \( T_i \) = initial average concentrate temperature
- \( k \) = concentrate thermal conductivity
- \( T_s \) = can surface temperature
- \( T_\infty \) = air temperature
An analytical solution to this system is prevented by the convective boundary condition.

The numerical solution employed an implicit technique using finite differences. In this case the governing difference equations were [11]

\[
\sum_{j} \frac{T_j^P - T_i^P}{R_{ij}} = C_i \frac{T_i^{P+1} - T_i^P}{\Delta t}
\]

(11)

where

- \( T_i^P \) = nodal temperature at time level \( P \)
- \( i \) = nodal location
- \( j \) = refers to each adjacent node
- \( C_i \) = lumped system heat capacitance for node \( i \)
- \( R_{ij} \) = thermal resistance between nodes \( i \) and \( j \)
- \( \Delta t \) = time step

The resistances and capacitances are calculated by

\[
C_i = \rho c \Delta V_i
\]

\[
R_{ij} = \frac{\Delta X_{ij}}{kA} \text{ for conduction}
\]

\[
R = \frac{1}{hA} \text{ for convection}
\]

where

- \( \rho \) = density
- \( c \) = specific heat capacity
- \( \Delta V_i \) = volume of \( i \)th element
- \( \Delta X_{ij} \) = distance between nodes \( i \) and \( j \)
- \( A \) = nodal area for heat transfer
Since an implicit method was used, the time step had to be chosen to meet adequate stability criteria. Discontinuities in the effective thermal properties around the initial freezing point required a small time step to assure a stable solution. A time step of 3.6 seconds was chosen for the 12 ounce can size.

Each cylinder of concentrate was divided into 3 sets of 3 concentrate rings for a total of 9 volume elements and nodes. Various volume element arrangements were considered. The arrangement used is sketched in figure 2. The width of the outermost elements, in either the axial or radial direction, is half that of the inner elements. This arrangement improved the stability of the solution over the case where nodes are spaced equally. Also, the outer elements are thinner and provide a closer approximation of the surface temperature.

In actual concentrate cans, a small air gap exists at the top of the can. The air gap tends to insulate the top surface of the concentrate. The size of this air gap was measured, and its thermal resistance was calculated. Since the air gap's thermal resistance is in series with the convective thermal resistance of the top surface, they were summed and used as an effective convective thermal resistance for the top surface.

An additional concern was that the air temperature changes as the air flows between the concentrate cans. This effect was accounted for by using equation (7) to calculate the heat transferred to the air, $q_{air}$. The temperature rise of the air, can be calculated from the definition of specific heat capacity and is given by
Fig. 2. Volume Element and Nodal Arrangement

Side view with can upright

End view
\[ \Delta T_a = \frac{q_{air}}{(c_p \dot{m})} \]

where

\[ \Delta T_a = \text{temperature rise of the air} \]

\[ c_p = \text{air specific heat capacity} \]

\[ \dot{m} = \text{mass flow rate of air} \]

Substituting equation (7) for \( q_{air} \) results in

\[ \Delta T_a = \frac{(h a' \dot{V} \Delta T_{In})}{(c_p \dot{m})} \]  

(12)

A value for the can's surface temperature \( T_s' \) is required for \( \Delta T_{In} \). The value of \( T_s \) was assumed to be uniform over each can's surface. The computer program estimated \( T_s \) by averaging the temperature of the outer elements, weighted relative to their surface areas, at each time step. The air temperature near the surface of each volume element, for use in equation (11), was then estimated by assuming the air temperature for the surfaces of the upstream volume elements was equal to the initial air temperature. The air temperature near the surface of the downstream volume elements was assumed to be equal to the initial air temperature plus the temperature rise calculated from equation (12). The air temperature used for the middle surfaces was the average of the air temperature used on the ends. The results of these assumptions agreed well with experimental observations.

A 16 element model was programmed, but its solution for average concentrate temperature varied only about 1°F from the 9 element model after a 30°F temperature change. Also, the time step needed for stability did not change. The 9 element model was chosen for the parameter study because it used about 25% less computer time.
The final model could predict the temperature distribution in cans of concentrate versus time in the tunnel. Time in the tunnel is related to position in the tunnel by the tunnel length and the average conveyor belt speed. The model was used primarily to predict concentrate freezing times for various values of upstream air temperature $T_a$, initial concentrate temperature, $T_i$, convective heat transfer coefficient, $h$, can height and radius, and concentrate thermal properties.
CHAPTER 3
MEASUREMENTS
3.1 MEASUREMENTS

A variety of measurements were necessary to evaluate the accuracy of the computer model, and to determine the tunnel's typical operating conditions. Measurements of the tunnel's internal operating conditions were complicated by the harsh environment created inside the tunnel. Also, the concentrate can size and average conveyor belt speed varied with production requirements. To simplify measurement problems, data was only recorded for the 12 ounce can size, which was the most frequent size cooled in the tunnel.

3.2 PRODUCTION RATE

The production rate, considered here as the mass flow rate of concentrate through the tunnel, depends on the average conveyor belt speed and the void fraction, $\varepsilon$. Although the conveyor belt drive was a constant speed drive, it was occasionally turned off and on due to production requirements. An average conveyor belt speed was estimated by noting the time required for a given can to go from entrance to exit of the tunnel. The average speed varied between 60 and 80 ft/hr.

The void fraction was estimated by using installed counters. Immediately after exiting the freeze tunnel, the cans were packed in boxes. Installed counters displayed the number of boxes that had been produced. The number of cans exiting the tunnel during the time required for a given can to pass from entrance to exit of the tunnel, was calculated from the counter readings. The bed volume was assumed to be
one can height tall, and as long and as wide as the conveyor belt inside the tunnel. Since, the mass and volume per can was chosen, the $\varepsilon$ could be estimated as

$$\varepsilon = 1 - \frac{(#\text{cans per bed})(\text{volume per can})}{(\text{bed volume})}$$

It was found that $\varepsilon$ typically varied between 0.4 and 0.5. An average value of 0.45 was estimated for the parameter study.

3.3 AIR FLOW RATE

It is necessary to determine the characteristic air velocity in the packed bed to predict a convective heat transfer coefficient. Figure 3 is a simple sketch of the tunnel air flow. Cold air is blown by 10 fans operating in parallel, through the mesh conveyor belt and the bed of concentrate cans. The air then flows through 8 evaporators operating in parallel, and returns to the fan suction. The fans were not spaced evenly along the length of the tunnel and the air velocities in the bed were higher near the ends of the tunnel than near the middle.

The average volumetric flow rate of air through each fan, $Q_f$, was estimated. Air velocities approaching 100 mph with air temperatures of about $-20^\circ\text{F}$ precluded involved or time consuming measurements in the vicinity of the fans. A pitot-static tube and an inclined oil manometer were used to measure the radial velocity distributions in the fan suction. It would have been more desirable to work on the discharge side of the fans, for safety reasons, but the fan discharge was not accessible during freeze tunnel operation due to the tunnel construction. Data were obtained for values of velocity and radial location along horizontal and vertical radials of several fan suction. Data
Fig. 3. Freeze Tunnel Air Flow Sketch
for at least 4 values of velocity and radial position were recorded for each radial considered.

The velocities measured were graphically integrated over the cross-sectional area of a fan suction to determine the volumetric flow rate per fan [12]

\[ Q_f = \int u dA \]

The average flow rate per fan was approximately 24,000 cfm. By assuming uniform flow, the characteristic velocity, \( u^* \), of the packed bed can be estimated for equation (10).

\[ u^* = \frac{Q}{\varepsilon A} \]

Then, \( u^* \) would be approximately 6 ft/sec.

Attempts were also made to measure the velocity distribution of the bed by directly measuring velocity in the void spaces, over the cross-sectional area of the bed. The manometer could not be used because the bed was in motion, and no level surfaces existed to put it on. A styrofoam ball type of flow detector was used with some success. Although the lower air velocities, in the larger void spaces, were below the detector's minimum sensitivity, it would consistently indicate the air velocities in the void spaces in the tightly packed regions of the bed. When averaged over the length of the tunnel, and corrected for temperature, the peak air velocity was approximately 9 ft/sec. This, of course, is not \( u^* \) but can be used to approximate its value.

When observed from above, the packed bed appears to consist of regions of tightly packed cans separated by small, relatively empty gaps. This observation suggested a way to use the peak air velocity to pre-
dict $u^*$. The bed is considered to consist of two types of areas, one of tightly packed cans and the other of no cans at all, such that

$$A_{\text{void}} = A_1 + A_2$$

where

- subscript $1$ = refers to the tight packed region
- subscript $2$ = refers to the region of no cans

Then equation (9) becomes

$$u^* = \frac{1}{A_{\text{void}}} \int u_1 dA_1 + \frac{1}{A_{\text{void}}} \int u_2 dA_2$$  \hfill (13)

The values of $A_1$ and $A_2$ can be estimated from $\varepsilon$ data. As previously discussed, on the average, $\varepsilon = 0.45$ for the tunnel. In the open regions, $\varepsilon_2 = 1.0$ by definition. The value of $\varepsilon$ in the tight packed region can be estimated from the tightest observed packing geometry as viewed from directly above the bed. Neglecting the edges of the region, every void space is surrounded by 3 cans and every can is surrounded by 6 void spaces. By observation the smallest unit of area that is characterized by a void fraction typical of the region, would be a triangular region, as sketched in figure 4. The length of each side is equal to twice the radius of a can. The equilateral triangle is drawn between the centers of any three adjacent cans. The minimum void fraction, $\varepsilon_1$, expected can be analytically or graphically estimated and is approximately 0.09. Of course, when $\varepsilon_1$ and $\varepsilon_2$ are averaged over the area of the bed, the average $\varepsilon$ must be 0.45 as previously determined,

$$\varepsilon = \frac{\varepsilon_1 A_1 + \varepsilon_2 A_2}{A_{\text{void}}}$$
Fig. 4. Top View of Observed Packing Arrangement
Also

\[
\frac{(A_1 + A_2)}{A_{\text{void}}} = l
\]

Combining the last two equations, eliminating \( A_2 \), setting \( \varepsilon_2 = 1.0 \) and solving the \( A_1 \) results in

\[
\frac{A_1}{A_{\text{void}}} = \frac{(1-\varepsilon)/(1-\varepsilon_1)}{(1-\varepsilon)/(1-\varepsilon_1)}
\]

Finally, equation (13) can be used to estimate \( u^* \). Values for \( A_1/A_{\text{void}} \) and \( A_2/A_{\text{void}} \) are determined from the equation and \( \varepsilon \) data above. A value for \( u_1 \) was measured. But the value of \( u_2 \) was below the minimum detectable velocity for the detector used. The temperature corrected minimum detectable velocity was approximately 1.3 ft/sec. When \( u_2 \) is assumed to have a value between 0.0 and 1.3 ft/sec, a value for \( u^* \) between 5.5 and 6.0 ft/sec results respectively. This result agrees with the value of 6.0 ft/sec resulting from the fan data.

3.4 TEMPERATURES

The average concentrate temperature was measured as a function of time and position in the freeze tunnel. Also, the air temperature upstream and downstream of the concentrate cans was measured as a function of position in the tunnel. These temperatures were measured with laboratory grade or precision grade mercury thermometers. Either partial or total immersion thermometers were used, as required by the measurement.

The steady state air temperatures were relatively consistent. The air temperature averaged -20°F upstream of the concentrate. The downstream air temperature varied with position in the tunnel. Near the entrance of the tunnel, the downstream air averaged -2°F, while at the exit the downstream air temperature averaged -15°F. However, necessary
Evaporator defrosting did temporarily affect the air temperatures. Evaporator defrosting occurred automatically for 1 evaporator every 3 hours. They were defrosted with hot gas. Hydraulically operated louvers were designed to automatically shut and isolate each evaporator during its defrost cycle and then open for normal operation. However, the louver system did not operate properly during the time period in which data was taken. The louvers remained open, or partially open, during defrost periods. Air temperatures downstream of a defrosting evaporator were observed to reach 30°F. This, of course, also affected the concentrate temperatures. Although it was attempted, taking data during defrosting periods could not be avoided because of the volume of data needed to establish typical operating conditions. Also, it usually took between 2.0 and 2.5 hours for a can of concentrate to go from tunnel entrance to exit so that most cans were subjected to a defrost cycle, which occurred every 3 hours.

Measuring the average temperature of a cylindrical concentrate can in a freeze tunnel is difficult. The major difficulty is caused by the large temperature gradient that results from the rapid freezing process. In some cases a can of concentrate, partway through the tunnel may be frozen solid near its surface, and still be liquid in the middle. Two methods were used to approximate the typical average temperature of the concentrate versus time and distance in the tunnel.

One method used to approximate the concentrate temperature, referred to as mixing cup method, was to empty selected cans into prechilled thermos bottles. The concentrate was then mechanically mixed until its
temperature was uniform enough to be measured with a single thermometer. However, the frozen concentrate was usually too hard to be easily mixed. Taking too much time or expending too much work mixing the concentrate was found to affect the concentrate temperature. A standard routine was established to expeditiously mix the concentrate. The routine sometimes left temperature variations within the concentrate of about 2F, but further mixing could also produce a comparable variation in the temperature. This measurement uncertainty contributed to some of the data scatter, primarily in the well frozen cans that had been in the tunnel over an hour. Data were collected by removing cans from specific locations in the tunnel and recording their temperatures. The average conveyor belt speed was measured and used to estimate the time the cans had been in the tunnel based on their positions. Data were collected several times on different days so that typical values could be determined. The actual data points obtained are plotted in figure 5 versus time in the freeze tunnel. The plot is dimensionless with the dimensionless temperature, \( \Theta \), defined as

\[
\Theta = \frac{T_i - T_a}{T_i - T_a}
\]

where

\( T_i \) = the average initial concentrate temperature

\( T_a \) = the air temperature upstream of the concentrate and the dimensionless time, \( \bar{\xi} \), defined as

\( \bar{\xi} = t/t_0 \)
Fig. 5. Temperature Data

* numeral next to data point indicates number of points at that location
where
\[ t_0 = \text{the reference time} \]

In this case
\[ T_i = 28^\circ F \]
\[ T_a = -20^\circ F \]
\[ t_0 = 2.5 \text{ hr} \]

The data scatter is due to a variety of reasons. Variations in initial concentrate temperature, measurement uncertainties, evaporators defrosting at different locations, variation in air flow rates between different regions of the tunnel and the stop and go operation of the conveyor belt are all contributors to the data scatter. However, the temperatures do generally decrease as expected. The average value of these data points is graphed versus time in the tunnel in figure 6. The average concentrate temperature decreases more slowly in the middle of the tunnel than near the ends. This is expected because of the higher effective specific heat capacity of the concentrate at temperatures typical of those in the middle of the tunnel and also because of the higher air velocities near the ends of the tunnel.

A different method for determining the average concentrate temperature was also used and is referred to as the computer aided method for discussion purposes. A hole was punched in the center of several can tops at the tunnel entrance. Mercury thermometers were then inserted in the cans of concentrate. Washers were taped to the thermometers to hold them at the proper immersion depth. The
Fig. 6. Average Temperature of Concentrate vs. Time in the Freeze Tunnel
temperatures indicated, and time of the readings, were recorded periodically as the cans progressed through the tunnel. The resulting data is also plotted in figure 5. Each data point is an average of the several thermometer readings recorded each time. These temperatures are not the average concentrate temperatures, but instead, the concentrate temperature at the can's centerline, near the thermometer's mercury bulb.

These thermometer readings were used to estimate the average concentrate temperature with the aid of the computer model. The computer could predict the average temperature and the temperature at 9 nodal locations in a can of concentrate as a function of time for any value of heat transfer coefficient, air temperature, initial concentrate temperature, and can size. All parameters of the computer model were set to the best estimated conditions in the tunnel. Values for the average temperature and the temperatures at the 3 centerline nodal positions were determined as function of time with the computer. Temperatures at the centerline nodal positions were used to obtain approximate graphs of temperature versus height at the can's centerline. The graphs were used to average the centerline temperature over the heights occupied by the mercury bulb. Comparing the predicted results with the measured results shows an average difference of less than $2.0^\circ F$ between the computer prediction and the measured center line temperatures for 2 hours of cooling. It was assumed the difference between the centerline temperature around the mercury bulb, predicted from the computer results, and the
computer predicted average temperature for the can was equal to
the difference between the measured centerline temperature and
the actual average temperature. By adding this temperature differ­
ence to the measured centerline temperature an estimate of the
corresponding average temperature was obtained. The temperature
difference varied with time so the procedure was repeated for
different times. A plot of the results of this computer aided
method is contained in figure 6, along with the results of the
mixing cup method. The dimensionless values are defined as in the
mixing cup method, except that the average initial concentrate
temperature was measured as 25°F rather than 28°F.

Comparing the graphs in figure 6 shows close agreement
except at times near the end of the tunnel. The computer aided
graph is longer because the average conveyor belt speed was slower
when that data was recorded, and the cans were in the tunnel longer.
The largest temperature difference between methods occurs at the
end of the mixing cup curve, when the cans were near the tunnel
exit. During periodic checks of the temperature of concentrate
exiting the tunnel temperature differences this large were observed
as a result of the routine operation of the tunnel. However, another
possible factor in this discrepancy is that when data was recorded
for the mixing cup curve the freeze tunnel door was open longer, as
thermos bottles were passed in and out, than when thermos bottles
were not used and only a data taker went in and out. The freeze
tunnel door was large and when open could significantly increase
the cooling load resulting in generally higher temperatures.

A final note on concentrate temperatures tends to agree equally with the results of both the mixing cup method and the computer aided method. The operators of the freeze tunnel set up its operation to produce a nominal concentrate outlet temperature of 0.0°F. Concentrate is normally stored in the 0.0°F to -5.0°F temperature range.

3.5 AVERAGE NUSSELT NUMBER

One initial use of the measurements is to evaluate the accuracy of and improve the precision of the computer model. Many of the parameters in the model, such as $u^*$, the packed bed characteristic velocity and the typical void fraction, could only be obtained by measurements.

The typical heat transfer coefficient and Nusselt number for the freeze tunnel can now be estimated based on the correlation in Section 2.3. For the 12 ounce can size (0.104 ft in radius and 0.375 ft in height), measured air temperatures, a characteristic velocity of 6 ft/sec, and a void fraction of 0.45, the correlation predicts a heat transfer coefficient and Nusselt number of about

$$ h = 8.2 \text{ Btu/hr ft}^2\text{°F} $$

$$ Nu = 126.0 $$

The accuracy of the correlation is better than $\pm 25\%$ [10]. Of course, the graphs of concentrate temperature versus time can also be used to measure $h$.

The computer is needed to estimate $h$ from the temperature
measurements. This can be done by simulating the existing conditions in the tunnel as closely as possible, and then varying h until the computer predicted temperatures are approximately equal to the measured temperatures. Simulating conditions in the tunnel required varying u* and h in different regions of the tunnel, to account for the uneven fan distribution, and also required simulating evaporator defrosting. Varying h was accomplished by dividing the tunnel into three regions, a region of low u* and h in the middle, and two regions of high u* and h at the ends of the tunnel. It was assumed u* was approximately 50% greater in the first 15% and the last 15% of the tunnel based on the measurements of the void space air velocities. To simulate the louver malfunction based on measured values, it was assumed that after 1 hour of cooling the upstream air temperature increased to 30°F for about 10 minutes and then returned to its original value of -20°F. The resulting predictions of average concentrate temperature versus time for various h and Nu is graphed in figure 7 along with the measured values.

Comparison of these curves shows the computer aided curve is very close to the curve for the average h and Nu predicted from the packed bed correlation. The mixing cup curve appears to be closer to a Nu of 86, even when some deviation is allowed to account for door openings. The average over the length of the tunnel of the two measured curves is close to the curve resulting from a Nu of 101. To be conservative, the tunnel's typical Nu was assumed to be 101 for the purposes of the parameter study. This value is
Fig. 7. Comparison of Measured and Computer Predicted Values of the Average Temperature of Concentrate vs. Time in the Freeze Tunnel
20% lower than the value predicted by the packed bed correlation. This accuracy is acceptable for an initial parameter study. A more accurate evaluation would require more temperature data or another method of evaluating the concentrate temperature.

3.6 IMPLICATIONS OF THE MEASUREMENTS ON THE PARAMETER STUDY

The measurements can be used to estimate the relative contributions of each factor in the tunnel energy balance. Any factor affecting the cooling load will affect the tunnel's coefficient of performance and the parameter study. Evaluation of each factor will help indicate the importance and potential of each factor to the efficient operation of the tunnel. Recalling from Section 2.1, the equations for the tunnel coefficient of performance and the tunnel energy balance are

\[ \text{COP}_t = \frac{q_c}{q_L/COP + W_f} \]  \hspace{1cm} (4)
\[ q_L = q_c + q_{\text{trans}} + q_{\text{inf}} + q_f \]  \hspace{1cm} (2)

where

- \( \text{COP}_t \) = tunnel coefficient of performance
- \( q_c \) = heat removal rate from the concentrate
- \( q_L \) = cooling load
- \( q_{\text{trans}} \) = transmission heat gain
- \( q_{\text{inf}} \) = infiltration heat gain
- \( q_f \) = fan heat gain

\( \text{COP} \) = coefficient of performance for the refrigeration plant

Calculation of the COP is discussed in Section 4.2. The other terms can be calculated using equations from Chapter 2, tunnel dimensions from
figure 1, and data from Chapter 3. A summary of the tunnel's typical operating conditions is contained in Table 2. Using this information, the energy consumption can be calculated.

The heat removal rate from the concentrate is calculated by

\[ q_c = \dot{m} \int_{T_i}^{T_f} c_{ef} \,dT \]  

(5)

When the integral is evaluated between 28°F and 0°F by graphically integrating the values in Table 2.1, the result is

\[ \int_{28}^{0} c_{ef} \,dT = 58.0 \text{ Btu/lbm} \]

The mass flow rate of the concentrate can be found using

\[ \dot{m} = S \, W \, H \,(1-c)\rho \]

where

\[ S = \text{conveyor belt speed} \]
\[ W = \text{packed bed or conveyor belt width} \]
\[ H = \text{packed bed height} \]
\[ \rho = \text{average concentrate density} \]

The average concentrate density, \( \rho \), was graphically averaged between 28°F and 0°F and is approximately 74 lbm/ft\(^3\). Using this information

\[ q_c = 6.20 \times 10^5 \text{ Btu/hr} \]

The transmission heat gain is calculated using data from Table 2. For the walls and ceiling the result is

\[ q = 2.35 \times 10^4 \text{ Btu/hr} \]

and for the floor the result is

\[ q = 2.40 \times 10^4 \text{ Btu/hr} \]
<table>
<thead>
<tr>
<th>Characteristic</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Tunnel Length</td>
<td>150 ft</td>
</tr>
<tr>
<td>Tunnel Width</td>
<td>20 ft</td>
</tr>
<tr>
<td>Tunnel Height</td>
<td>15 ft</td>
</tr>
<tr>
<td>Conveyor Belt Length</td>
<td>150 ft</td>
</tr>
<tr>
<td>Conveyor Belt Width</td>
<td>10 ft</td>
</tr>
<tr>
<td>Average Conveyor Belt Speed</td>
<td>70 ft/hr</td>
</tr>
<tr>
<td>Packed Bed (can) Height</td>
<td>0.375 ft</td>
</tr>
<tr>
<td>Void Fraction</td>
<td>0.45</td>
</tr>
<tr>
<td>Door Height</td>
<td>8 ft</td>
</tr>
<tr>
<td>Door Width</td>
<td>4 ft</td>
</tr>
<tr>
<td>Overall Heat Transfer Coefficient, Walls</td>
<td>0.025 BTU/hr ft(^{2})°F</td>
</tr>
<tr>
<td>Overall Heat Transfer Coefficient, Floor</td>
<td>0.1 BTU/hr ft(^{2})°F</td>
</tr>
<tr>
<td>Design Wet Bulb Temperature</td>
<td>79°F</td>
</tr>
<tr>
<td>Design Dry Bulb Temperature</td>
<td>93°F</td>
</tr>
<tr>
<td>Design Ground Temperature</td>
<td>60°F</td>
</tr>
<tr>
<td>Initial Concentrate Temperature</td>
<td>28°F</td>
</tr>
<tr>
<td>Final Concentrate Temperature</td>
<td>0°F</td>
</tr>
<tr>
<td>Freeze Tunnel Air Temperature</td>
<td>-20°F</td>
</tr>
</tbody>
</table>
The total transmission heat gain is

\[ q_{\text{trans}} = 4.75 \times 10^4 \text{ Btu/hr} \]

More assumptions are required to calculate \( q_{\text{inf}} \). The enthalpy difference between outside air at design temperature and inside air at \(-20^\circ\text{F}\), \( \Delta h \), is approximately 50 Btu/lbm. The cfm infiltrating with the cans is

\[ \text{cfm} = S WH e \]

Using values from table 2 results in a heat gain from infiltration with the conveyor belt of

\[ q = 98 \text{ Btu/hr} \]

The infiltration heat gain due to door openings can vary widely. Assuming the door is open only 15 seconds per hour on the average and using the methods discussed in Section 2.1 results in a heat gain due to door openings of

\[ q = 4.50 \times 10^3 \text{ Btu/hr} \]

The total infiltration heat gain is

\[ q_{\text{inf}} = 4.60 \times 10^3 \text{ Btu/hr} \]

The heat gain resulting from the fans is

\[ q_f = 3.00 \times 10^5 \text{ Btu/hr} \]

The total cooling load from equation (2) is

\[ q_L = 9.72 \times 10^5 \text{ Btu/hr} \]

Table 3 is a summary of these results.

The major loads are \( q_c \) and \( q_f \) while the value of \( q_{\text{inf}} \) is negligible. The value of \( q_{\text{trans}} \) is relatively small even as a worst case value. Therefore, for the purposes of the parameter study the cooling load is
### TABLE 3
SUMMARY OF COOLING LOADS

<table>
<thead>
<tr>
<th>Source</th>
<th>Heat Gain (BTU/hr)</th>
<th>Percent of Total</th>
</tr>
</thead>
<tbody>
<tr>
<td>Concentrate</td>
<td>6.20 x 10^5</td>
<td>63.8</td>
</tr>
<tr>
<td>Transmission</td>
<td>4.75 x 10^4</td>
<td>4.9</td>
</tr>
<tr>
<td>Infiltration</td>
<td>4.60 x 10^3</td>
<td>0.5</td>
</tr>
<tr>
<td>Fans</td>
<td>3.00 x 10^5</td>
<td>30.9</td>
</tr>
<tr>
<td>Total</td>
<td>9.72 x 10^5</td>
<td>100.0</td>
</tr>
</tbody>
</table>
approximated as

\[ q_L = q_c + q_f \]

Then, equation (4) becomes

\[ \text{COP}_t = \frac{q_c}{(q_c + q_f)/\text{COP} + W_f} \]  \hspace{1cm} (14)
CHAPTER 4
THE PARAMETER STUDY
4.1 THE COEFFICIENT OF PERFORMANCE OF THE FREEZE TUNNEL

The equation for $\text{COP}_t$ can still be put in a more convenient form for the parameter study. Rearranging the terms in equation (14) results in

$$\text{COP}_t = \frac{\text{COP}}{1 + (q_f + W_f \text{COP})/q_c}$$

From Chapter 2,

$$q_f = 2995 \text{ hp}$$

where

$$W_f = \text{fan horsepower}$$

Also, using a convenient conversion results in

$$W_f = 2545 \text{ Hp}$$

Substituting these expressions into the $\text{COP}_t$ formula results in

$$\text{COP}_t = \frac{\text{COP}}{1 + (2995 + 2545 \text{ COP}) \frac{H_p}{q_c}}$$

One more simplifying approximation is to assume this equation can be rewritten

$$\text{COP}_t = \frac{\text{COP}}{1 + (1+\text{COP})K}$$

In this form, $K$ is a dimensionless number defined as

$$K = \frac{2545 \text{ Hp} \Delta t}{\Delta Q_c}$$

where

$$\Delta Q_c = \text{the nominal heat removed from the concentrate in a full freeze tunnel}$$

$$\Delta t = \text{the freezing time or the time required for a can to pass through the freeze tunnel}$$
and

\[ q_c = \frac{\Delta Q_c}{\Delta t} \]

\( K \) is a ratio of the work done by the fans divided by the useful refrigeration effect. Even before the parameter study, it is easy to see the significance of this ratio. For a given COP, the tunnel is most effective, or the COP\(_t\) is a maximum, when

\[ K = 0 \]

or

\[ q_c \gg \dot{W}_f \]

Of course when this occurs, there are essentially no fans in the tunnel and the freeze tunnel has become a refrigerated space. Typically this cannot be accomplished because the freezing times become too long. To obtain the desired freezing times, for a given capacity of the tunnel, fans are added. As K increases, the COP\(_t\) decreases. Figure 8 and figure 9 show briefly how COP\(_t\), COP, and K are related.

4.2 VARIATION IN THE COEFFICIENT OF PERFORMANCE OF THE REFRIGERATION PLANT WITH FREEZE TUNNEL AIR TEMPERATURE

The COP for the refrigeration plant associated with the observed freeze tunnel is difficult to accurately calculate. The two stage ammonia vapor compression plant supplies loads other than the freeze tunnel. Some loads are supplied from the intermediate stage. The enthalpy of the refrigerant cannot be estimated for all the important thermodynamic states. For the purposes of the parameter study, the COP was estimated by assuming the refrigerant reaches each compressor as a saturated vapor, compression is isentropic, and the
Fig. 8. Freeze Tunnel Coefficient of Performance Divided by the Refrigeration Coefficient of Performance vs. Fan Power Divided by the Useful Refrigeration Effect
Fig. 9. Freeze Tunnel Coefficient of Performance vs. Refrigeration Unit Coefficient of Performance
minimum enthalpy at any pressure is approximately equal to the enthalpy of saturated liquid refrigerant at the highest pressure in the cycle.

Information has to be obtained concerning operating pressures to calculate the COP to be used in the parameter study. The stage pressures that produce an air temperature in the tunnel of -20°F were observed as

<table>
<thead>
<tr>
<th>Pressure Type</th>
<th>Pressure Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>High Pressure</td>
<td>170 psig</td>
</tr>
<tr>
<td>Intermediate Pressure</td>
<td>30 psig</td>
</tr>
<tr>
<td>Evaporator Pressure</td>
<td>10 inch Hg, vac</td>
</tr>
</tbody>
</table>

The air temperature is a parameter in the study. To change the air temperature the evaporator pressure must be changed, for a given cooling load. This of course affects the COP, so that at every air temperature and evaporator pressure, a new COP must be estimated. As a rough estimate, it was assumed that for a given change in air temperature, the evaporator's saturation temperature must change an equal amount. It was also assumed that the evaporator pressure changes a corresponding amount while the high and intermediate pressures are constant. Using these assumptions, the evaporator pressures needed to produce given air temperatures and the corresponding COP's were estimated. The results are listed in table 4.

4.3 VARIATION IN THE FAN WORK WITH NUSSELT NUMBER

Fan work, $W_f$, is also an important parameter. $W_f$ is related to the packed bed's characteristic velocity, $u^*$, and it's heat transfer coefficient, $h$. However, predicting how a change in $W_f$ will affect $u^*$ and $h$ is a difficult problem. Usually detailed know-
<table>
<thead>
<tr>
<th>COP</th>
<th>Air Temperature (°F)</th>
</tr>
</thead>
<tbody>
<tr>
<td>6.0</td>
<td>-10</td>
</tr>
<tr>
<td>5.4</td>
<td>-20</td>
</tr>
<tr>
<td>4.9</td>
<td>-30</td>
</tr>
<tr>
<td>4.5</td>
<td>-40</td>
</tr>
</tbody>
</table>
ledge of the fan's characteristic curve of pressure head versus volumetric air flow rate is needed, as well as the systems characteristic curve of head loss versus volumetric air flow rate, to accurately estimate a change in a system's operating point [13]. In this case, only 1 operating point is known. The system curve could be determined experimentally, but this would be too difficult. Therefore, for an initial investigation the fan laws [13] were used as a rough approximation for the relationship between \( W_f \) and \( u^* \). The applicable fan law in this case, assuming \( u^* \) is directly proportional to the volumetric air flow rate, is

\[
W_f = u^*^3
\]

At the known operating point,

\[
W_f = 100 \text{ Hp} \\
u^* = 6.0 \text{ ft/sec} \\
Nu = 101
\]

For a given change in \( W_f \), the fan law can be used to estimate the new \( u^* \). Then the packed bed correlation can be used to calculate the new \( Nu \) and \( h \).

4.4 THE FREEZE TUNNEL COEFFICIENT OF PERFORMANCE VERSUS THE NUSSELT NUMBER

In conducting a parameter study for the freeze tunnel under observation, a simplification occurs because \( \Delta Q_c \) is fixed. Any change in \( Nu \), with a corresponding change in \( W_f \), affects both the COP\(_t\) and the freezing time, \( \Delta t \). The computer model can be used to predict how a change in \( Nu \), or a change in air temperature, \( T_a \), will affect the freezing time. Although the \( \Delta Q_c \) is fixed, any change in the freezing
time will affect $q_c$.

The computer was programmed to predict the average concentrate temperature versus time for various values of $h$ and $T_a$. The time required for the average concentrate temperature to change from an initial value of $28^\circ F$ to a final value of $0^\circ F$, considered the freezing time, was determined from the computer output. Then the $COP_t$ was calculated for each value of $h$ and $T_a$, or equivalently, $Nu$ and $T_a$. For every value of $T_a$, the estimated $COP$ from Table 4 was used to estimate $COP_t$. The fan horsepower, $Hp$, was estimated for each value of $h$ and $Nu$, by using the fan law discussed in Section 4.3, relative to the known operating conditions. Since $Hp$, $\Delta Q_c$, and $\Delta t$ are known at each point, $K$ may also be calculated. A summary of the results is graphed in Figure 10 and figure 11.

Figure 10 is a graph of $COP_t/COP$ versus $Nu$. The ratio of $COP_t/COP$ has a maximum value of 1.0. When $COP_t$ equals $COP$, the least energy is expended for a given useful refrigeration effect, $q_c$. Freeze tunnels are operated with lower efficiencies when it is necessary to provide a high $q_c$ and/or a short freezing time, $\Delta t$. When heat transfer is increased by using fans to increase the $Nu$ and the ratio of fan work divided by useful refrigeration effect $K$, increases, then $COP_t$ becomes less than $COP$. This relationship is displayed by equation (15) as well as figure 10.

Figure 11 is more informative because it shows more clearly how the relationship between $COP_t$ and $Nu$ is affected by $T_a$. At very low $Nu$, an increase in $T_a$ also increases the $COP_t$. This is because
Fig. 10. Freeze Tunnel Coefficient of Performance Divided by the Refrigeration Unit Coefficient of Performance vs. the Nusselt Number
Fig. 11. Freeze Tunnel Coefficient of Performance vs. the Nusselt Number
at low Nu, the system is closer to a refrigerated space than a freeze tunnel and the dominant effect of increasing $T_a$ is the corresponding increase in COP. But at high Nu, the dominant effect of an increase in $T_a$ is increased freezing time, and the COP$_t$ actually decreases. This seems to suggest that while maintaining a higher $T_a$ in a refrigerated space results in higher a COP and lower energy consumption in a refrigeration problem, in a freeze tunnel problem maintaining a higher $T_a$ results in a lower COP$_t$ and higher energy consumption. Also, the COP$_t$ decreases rapidly as $K$ increases, as expected.

4.5 THE FREEZE TUNNEL COEFFICIENT OF PERFORMANCE VERSUS FREEZE TUNNEL CAPACITY

The result of adding fans to a refrigerated space is to increase the rate of heat transfer. This increase in the rate of heat transfer increases the tunnels capacity, $q_c$, and for a given tunnel size decreases the freezing time. The price of the increased capacity is a decrease in COP$_t$. The relationship between COP$_t$ and $q_c$ for the observed tunnel is easy to determine at this point.

As a result of Section 4.4, values of COP$_t$, Nu, $T_a$, $K$, and freezing time, $\Delta t$, have already been estimated for a variety of computer simulated operating points. Since $\Delta Q_c$ is fixed, and a relationship between COP$_t$ and $\Delta t$ has been established, values of COP$_t$ versus $q_c$ can be generated from

$$q_c = \frac{\Delta Q_c}{\Delta t}$$

Figure 12 is a graph of COP$_t$ versus $q_c$ for the range of Nu
Fig. 12: Freeze Tunnel Coefficient of Performance vs. Freeze Tunnel Capacity
and $T_a$ investigated. The graph displays the important trends discovered in the previous graphs: the highest $COP_t$ is obtained for the lowest values of $K$ and $Nu$, and for high values of the $Nu$, the highest $COP_t$ is obtained for the lowest value of $T_a$. But it also shows that large values of $K$ restrict the tunnel to relatively low $COP_t$'s, for any value of $T_a$.

The highest $COP_t$'s exist at the lowest $Nu$ as expected. But relatively large capacities appear possible even for the lowest $Nu$ investigated. The $COP_t$ of the tunnel is high for low $Nu$ primarily because $K$ is so low. $K$ was calculated using the fan $Hp$ predicted by the fan laws [13]. For a $Nu = 46$, the fan law predicts a $Hp=2$ horsepower, relative to 100 horsepower for a $Nu=101$ as discussed in Section 4.3. Actually producing a significant cooling air flow in a freeze tunnel similar to the one observed with only 2 horsepower may not be achievable because of the physical size and flow characteristics of the evaporators and packed bed. Careful experimental analysis using system and fan curves [13] would be necessary to accurately predict behavior for any conditions significantly different from the measured conditions.

4.6 ENERGY COSTS VERSUS FREEZE TUNNEL CAPACITY

Considering the effects of various values of $Nu$, $T_a$ and $K$ on the $COP_t$ is important because the $COP_t$ is a measure of the tunnel's effectiveness. But a more obvious method of judging freeze tunnel performance is to consider its energy consumption per unit of processed food. The energy consumed, or equivalently the net work
expended, is related to $q_c$ by the definition of COP$\text{t}$ from equation (1)

$$W_c + W_f = q_c / \text{COP}_\text{t}$$

The monetary cost of the electricity to operate the tunnel is related to the work performed by

$$d = R \frac{q_c}{\text{COP}_\text{t}}$$

where

$$d = \text{hourly charge}$$

$$R = \text{cost per unit energy}$$

The unit monetary cost, or cost per unit of food processed is

$$D = d \Delta t = R \frac{Q}{\text{COP}_\text{t}}$$

where

$$D = \text{unit cost}$$

$$\Delta t = \text{freezing time}$$

$$Q = \text{heat removed per unit food product}$$

Figure 13 is a graph of hourly energy costs, and rate of energy consumption, versus capacity. Once again, except at low capacities, the least energy is expended for a given production rate at the lowest achievable values of $N_u$, $K$, and $T_a$. As the capacity is increased by lowering $T_a$ or increasing $N_u$ the costs increase. However, increased costs may be acceptable or even desirable if the increased capacity results in a decreased unit cost.

The unit costs, both in energy and money is graphed versus capacity in figure 14. This graph displays all the trends noted previously. The most useful new information displayed in this
Fig. 13. Freeze Tunnel Electricity Cost vs. Freeze Tunnel Capacity
Fig. 14. Freeze Tunnel Electricity Cost per Unit Processed vs. Freeze Tunnel Capacity
graph is that the ratio of the fan work divided by the tunnel capacity, $K$, can be related approximately to the unit cost of the product. Also, the unit cost and $K$ can be substantially reduced, for a given $Nu$ by lowering $T_a$.

4.7 RESULTS CONCERNING THE OBSERVED FREEZE TUNNEL

The parameter study applies directly to the observed freeze tunnel with the 12 ounce can size. However, the operating conditions measured in the tunnel reflected the malfunctioning defrosting louvers that resulted in less efficient operation than should occur nominally. The typical operating conditions from section 3.6 result in

$$K = 0.36$$

for an average $Nu$ of 101. The corresponding freezing time was about 1.9 hours. However, freezing times were observed to vary from approximately 1.7 hours up to 2.5 hours depending on day to day operating conditions. As a result $K$ varies from 0.33 to 0.48. This wide variation makes it difficult to predict the tunnel's COP$_t$ or energy costs at a given time. In this case, with a COP of 5.4 for a corresponding air temperature of $-20^\circ F$, the COP$_t$ is 1.33 for a $K$ of 0.33, and the COP$_t$ is 1.74 for a $K$ of 0.48. The variation in the COP$_t$ is about $\pm 11\%$ from its average value of 1.5.

When the defrost cycle is left out of the computer program, and the air temperature $T_a$ is constant at $-20^\circ F$, the predicted freezing time for a $Nu$ of 101 is about 1.6 hours. Then,

$$K = 0.31$$
COP_t = 1.8

For any of these operating points, figure 14 predicts a lower cost if the tunnel is operated with colder air temperatures and a smaller Nu, or less fans. For example, assume that the tunnel already operates at its highest expected efficiency with

\[ T_a = -20^\circ F \]
\[ Nu = 101 \]
\[ K = 0.31 \]
\[ COP_t = 1.8 \]

Figure 14 predicts for that operating condition a unit cost of approximately $0.80 per 100 cases. Each case contains 24 cans of the 12 fluid ounce size. Figure 14 also predicts that when

\[ T_a = -25^\circ F \]
\[ Nu = 90 \]

the unit cost is $0.55 per 100 cases, a savings of about 30%.

Using the fan law discussed in Section 4.3, only about 50 Hp in fans is required to produce \( Nu = 90 \). This is half the fans currently in the tunnel. Although the unit energy costs are small compared to the cost of the concentrate, it costs well over $2000.00 per month in electricity to run the fans, and to remove the heat they generate from the tunnel.

Unfortunately, operating the freeze tunnel is more complicated than assumed in this analysis. The concentrate inlet temperature varies, the can size varies, and the rate of production varies. Therefore, one optimum operating point cannot be chosen. Although
a Nu of 90 and $T_a$ of -25°F will adequately cool 12 ounce cans, it may not adequately cool a larger can unless the conveyor belt and production is slowed down. Some reserve capacity is needed. In general it would be desirable to maintain the lowest reasonable air temperature when the tunnel is operated at a higher capacity. If the product or cooling load changes, fans should be turned on and off as necessary to provide the desired exit temperature.

In the freeze tunnel observed, fans could not be secured selectively because all the fans operate in parallel with common inlet and outlet plenums. Some type of automatic damper system would be required to shut when the fan was secured to prevent reverse air flow through the idle fan. Assume a lower air temperature would allow 1 fan to be secured for half its normal operating time. During a 9 month season, almost $1000.00 could potentially be saved in electricity costs for that fan. Savings this large could justify an inexpensive damper system.

4.8 RESULTS FOR FREEZE TUNNEL DESIGN IN GENERAL

Constructing a freeze tunnel is one of the largest initial expenses when building a food processing plant [1]. It is obvious that minimizing $K$ will reduce the operating expenses of the tunnel. But the design must take into account trade-offs between the initial investment capital and final operating expenses. However, many important trends that apply to freezing 12 ounce cans of citrus concentrate will have some relevance to any freeze tunnel where the needed useful refrigeration effect needed is large.
The lowest operating costs occur at the coldest air temperatures because the freezing times are shorter, and the tunnel's capacity is greater. On the other hand, for a given air temperature, the lowest operating costs occur for the lowest $Nu$, and consequently the longest freezing time. This means the freeze tunnel should be designed to provide an adequate freezing time, but no shorter than necessary. The fans should be chosen to provide this freezing time with the lowest reasonably producable air temperature. If the capacity varies, the fans should be operated selectively to maintain the lowest effective value of $K$. 
CHAPTER 5

CONCLUSIONS
5.1 LIMITATIONS OF STUDY

Use of the fan laws to relate the Nusselt number and characteristic air velocity to fan horse power limits the range of accurate predictions to operating conditions close to those of the observed tunnel. Predictions for operating conditions significantly different from the observed conditions are only rough approximations. Many of the operating conditions investigated in theory may not be achievable in application.

For example, a minimum fan horsepower may be required to produce any significant air flow through the evaporators and food product. This limitation on the minimum fan horsepower was not considered in this report.

The investigation assumed that characteristics of the refrigeration plant, and associated COP and temperatures, were for a two-stage ammonia vapor compression plant. The COP's used were rough approximations. Different COP's and air temperatures may be achievable with different types of equipment.

Also, the relationship between the Nusselt number, the air temperature, and the freezing time varies with the product cooled,
and the range of temperatures through which the product is cooled. Because of this, the graphs and numerical estimates may not be applicable to freeze tunnels cooling significantly different products.

5.2 SIGNIFICANT RESULTS

The major value of this investigation is the trend and the relationship between, the energy consumption, the Nusselt number, the air temperature, and the ratio of fan work divided by useful refrigeration effect. In review, the most economical energy consumption occurs, for large freeze tunnels, when the freezing times are no shorter than required, the air temperature is the lowest achievable value, and the ratio of fan work divided by useful refrigeration effect is the lowest achievable value.

Although some efficiency of the refrigeration unit is lost by producing a low air temperature, this trend is more than offset by the increased freeze tunnel capacity, and freeze tunnel coefficient of performance.

The significance of the ratio of fan work divided by the useful refrigeration effect, K, was also important. The energy expended to produce the desired cooling effect per unit of food product is determined predominantly by the value of K for the freeze tunnel. The minimum achievable values of K depend on the freeze tunnel design, the required refrigeration effect, and properties of the food product. The range of values for K measured for the observed freeze tunnel are reasonably accurate and could
be used to compare the effectiveness of the observed tunnel to another.

The equation for the freeze tunnel's coefficient of performance, equation (15),

$$\text{COP}_t = \frac{\text{COP}}{1 + (1+\text{COP})K}$$

can be used by freeze tunnel operators and designers to estimate the freeze tunnel effectiveness. The heat content of the food product and the amount of food product in the tunnel would have to be determined. The freezing times can be estimated for a variety of products, and should be known by the tunnel operator. The horsepower of the fans is fixed, or determinable, so $K$ may be calculated frequently without any other knowledge than that of the freeze tunnel design and the thermal properties of the food product. If the COP can be estimated, then $\text{COP}_t$ can also be estimated.

5.3 POSSIBILITIES FOR FURTHER RESEARCH

It would be interesting to estimate and collect operating conditions for as wide a variety of freeze tunnels as available. One value of this would be to determine what minimum values of $K$ may be achieved for specific food products. This information would be valuable in minimizing energy consumption in future designs or modifications to existing equipment.

Another valuable result of finding more operating points is that figures similar to those in this report could be generated with more accuracy, and potentially used as standards or guides for freeze tunnel design. Fan laws and other simple approximations
could be used to generate the portions of the graphs between the known operating points.
APPENDIX:

SAMPLE COMPUTER LISTING
$JOB
C CALCULATION OF AVERAGE TEMPERATURE OF A PACKED BED OF JUICE CONCENTRATE CANS IN A FREEZE TUNNEL FOR VARIABLE THERMAL PROPERTIES AND A NINE ELEMENT MODEL
DIMENSION T (3,3), TP(3,3), V(3,3), V(3,3), H(2), K(3,3), S(3,3), HC(3,3)
DIMENSION C (3,3), K1 (3,3), K2(3,3)
REAL K, K1, K2, KI, LMTD
DO 210 NNT = 1, 4
READ (5,11) H(1), H(2)
11 FORMAT (2F10.0)
READ (5,12) TI, TA1, TA2, TA3
12 FORMAT (4F10.0)
READ (5,13) HGT, RAD
13 FORMAT (2F10.0)
READ (5,14) U, TL, DELT
14 FORMAT (3F10.0)
C DEFINE AREAS
AEND = 3.1415926 * RAD ** 2
ACYL3 = 6.2831853 * RAD * HGT
ACYL2 = 0.8 * ACYL3
ACYL1 = 0.4 * ACYL3
C DEFINE VOLUMES
VTOT = HGT * AEND
V(1,0) = 0.04 * VTOT
V(1,2) = 0.12 * VTOT
V(1,3) = 0.09 * VTOT
V(2,1) = 2.0 * V(1,1)
V(2,2) = 2.0 * V(1,2)
V(2,3) = 2.0 * V(1,3)
V(3,1) = V(1,1)
V(3,2) = V(1,2)
V(3,3) = V(1,3)
C DEFINE INITIAL TEMPERATURE DISTRIBUTION AND THERMAL PROPERTIES
   SI = 75.2
   KI = 0.18
   HCI = 0.73
   DO 30 I = 1,3
   DO 40 J = 1,3
   S(I,J) = SI
   K(I,J) = KI
   HC(I,J) = HCI
   C(I,J) = S(I,J) * HC(I,J) * V(I,J)
   T(I,J) = TI
40  CONTINUE
30  CONTINUE
C DEFINE MISC TERMS
   DELZ = HGT / 2.0
   DELR = RAD / 2.5
   TAVE = TI
   TA = TA1
   DI = 100.0
   DT1 = DELT * 60.0 / DT
   DT2 = DELT * 60.0
   ER = 0.0
   L = 1
   M = 1
   LO = 0
C DEFINE RESISTANCES
C AXIAL CONDUCTION
   AKEND = (0.75 * DELZ) / (AEND)
   AK1Z = AKEND / 0.16
   AK2Z = AKEND / 0.48
   AK3Z = AKEND / 0.36
C RADIAL CONDUCTION
   AK11R = (DELR * 6.0) / ACYL1
   AK12R = (DELR * 3.0) / ACYL2
AK21R = 0.5 * AK11R
AK22R = 0.5 * AK12R
WRITE (6,90) H(1), TI, RAD, TL, H(2), TA, HGT, U, DT1, DT2
90 FORMAT ('1', 'VARIABLE THERMAL PROPERTIES FOR BRIX NO 44.8' // ' H(1)
1=','F5.2,' BTU/HR-FT2-F',4X,'TI =', F6.1, ' F',5X,'RAD =', F7.4, 'FT',
25X,'TL =', F6.1, ' FT' // ' H(2) =', F5.2,' BTU/HR-FT2-F',4X,'TA =',
3F6.1, ' F',5X, HGT = ', F7.4,' FT',5X,' U =', F6.1,' FT/HR' // '1x,F5.3,
4' MINUTE ITERATIONS AND DATA IS PRINTED EVERY ',F5.3,' MINUTES'
5///)
WRITE (6,99)
99 FORMAT (1X, 5H TIME, 10H T(1,0) , 10H T(2,1) , 10H T(3,1) , 110H T(1,2) , 10H T(2,2) , 10H T(3,2) , 10H T(1,3) , 10H T(2,3) , 10H TAVE , 10H TA3 //)
WRITE (6,100) LO, T, TAVE, TA3
100 FORMAT (16, 11F10.1)
C COMPUTE NEW TEMPERATURE DISTRIBUTIONS
DELT = DELT / DT
DO 70L = 1, 100
C DEFINE FLOW DEPENDANT RESISTANCES
C END TO AMBIENT
RHEND = 1.0 / (H(M) * ΔEND)
RH11Z = RHEND / 0.16
RH12Z = RHEND / 0.48
RH13Z = RHEND / 0.36
C SIDES TO AMBIENT
RH13R = 1.0 / (H(M) * ACYL3)
RH13R = 4.0 * RH13R
RH23R = 2.0 * RH13R
DO 75 N = 1, 100
TP (1,1) = DELT/C(1,1)*(TA1/RH11Z+K(1,1)*(T(1,2)/AK11R+T(2,1)/AK1Z
1+1.-DELT/C(1,1)*1./RH11Z+K(1,1)*1./AK11R+1./AK1Z))*T(1,1)
TP(1,2) = DELT/C(1,2)*(TA1/RH12Z+K(1,2)*(T(1,1)/AK11R+T(1,3)/AK12R
1+T(2,2)/AK2Z))+1.-DELT/C(1,2)*1./RH12Z+K(1,2)*1./AK11R+1./AK12R
2+1./AK2Z))*T(1,2)
TP(1,3) = DELT/C(1,3)*(TA1/RH13R+TA1/RH13Z+K(1,3)*(T(1,2)/AK12R+T(1,2)/AK13R+1.)}
\[ \frac{12,3}{AK3Z}) + (1 - \frac{DELT}{C(1,3)} \right) \left( \frac{1}{RH13Z \cdot 1} + \frac{1}{RH13R} + \frac{K(1,3)}{1.AK12R+1.2/AK3Z}) \right) \] 

\[ TP(2,1) = \frac{DELT}{C(2,1)} \cdot K(2,1) \left( \frac{T(1,1)}{AK12R} + \frac{T(2,2)}{AK21R} + \frac{T(3,1)}{AK12Z} \right) + (1 - \frac{DELT}{C(2,1)} \right) \left( \frac{1}{RH23R} + \frac{1}{AK23R+1.1AK23Z}) \right) \] 

\[ TP(2,2) = \frac{DELT}{C(2,2)} \cdot K(2,2) \left( \frac{T(2,1)}{AK21R} + \frac{T(1,2)}{AK22R} + \frac{T(3,2)}{AK22Z} \right) + (1 - \frac{DELT}{C(2,2)} \right) \left( \frac{1}{RH23R} + \frac{1}{AK23R+1.1AK23Z}) \right) \] 

\[ TP(2,3) = \frac{DELT}{C(2,3)} \cdot \frac{TA2}{RH23R} \cdot (K(2,3) \left( \frac{T(2,2)}{AK22R} + \frac{T(1,3)}{AK3Z} \right) + (1 - \frac{DELT}{C(2,3)} \right) \left( \frac{1}{RH23R} + \frac{1}{AK23R+1.1AK23Z}) \right) \] 

\[ TP(3,1) = \frac{DELT}{C(3,1)} \cdot \frac{TA3}{RH13Z} \cdot (K(3,1) \left( \frac{T(3,2)}{AK12R} + \frac{T(1,1)}{AK12Z} \right) + (1 - \frac{DELT}{C(3,1)} \right) \left( \frac{1}{RH13Z} + \frac{1}{AK12Z+1.1AK12R}) \right) \] 

\[ TP(3,2) = \frac{DELT}{C(3,2)} \cdot \frac{TA3}{RH13R} \cdot (K(3,2) \left( \frac{T(3,1)}{AK12R} + \frac{T(2,3)}{AK3Z} \right) + (1 - \frac{DELT}{C(3,2)} \right) \left( \frac{1}{RH13R} + \frac{1}{AK12R+1.1AK12Z}) \right) \] 

\[ TP(3,3) = \frac{DELT}{C(3,3)} \cdot \frac{TA3}{RH13Z} \cdot (K(3,3) \left( \frac{T(3,2)}{AK12R} + \frac{T(2,3)}{AK3Z} \right) + (1 - \frac{DELT}{C(3,3)} \right) \left( \frac{1}{RH13Z} + \frac{1}{AK12Z+1.1AK12R}) \right) \] 

C REDEFINE TEMPERATURE DEPENDENT THERMAL PROPERTIES AND TEMPERATURES

DO 50 I = 1,3
DO 60 J = 1,3
T(I,J) = TP(I,J)
IF ( T(I,J) G T 15.5) GO TO 51
IF ( T(I,J) L -19.0) GO TO 52
K1(I,J) = 0.00074* T(I,J)**3 - 0.0040* T(I,J)**2 + 0.0032*

K2(I,J) = 0.00013* T(I,J)**3 + 0.0042* T(I,J)**2 + 0.017* T(I,J)

HC(I,J) = 0.0046* T(I,J)**2 + 0.14* T(I,J) + 2.0
IF ( T(I,J) L -19.0) HC(I,J) = 1.00
C(I,J) = S(I,J) * HC(I,J) * V(I,J)
GO TO 60

51 CONTINUE
K(I,J) = KI
HC(I,J) = HCI
S(I,J) = SI
C(I,J) = S(I,J) * HC(I,J) * V(I,J)
GO TO 60

52 CONTINUE
K(I,J) = 0.65
HC(I,J) = 1.00
S(I,J) = 71.77
C(I,J) = S(I,J) * HC(I,J) * V(I,J)

60 CONTINUE
50 CONTINUE

C CALCULATE RETURN AIR TEMPERATURE
ATOT = 2.0*AEND + ACYL3
TS = (T(2,3)*0.5*ACYL3+(T(1,1)+T(3,0))*0.16*AEND+(T(1,2)+T(3,2)*
10.48*AEND+(T(1,3)+T(3,3)))*(0.36(AEND+0.25*ACYL3))/ATOT
LMTD = (TA3 - TA1)/ALOG((TS - TA1)/(TS - TA3)
TA3 = TA1 + 0.24 * LMTD
TA2 = (TA1 + TA3)/2.0

75 CONTINUE

C CALCULATE AVERAGE TEMPERATURE
TAVE = 0.0
DO 80 I = 1,3
DO 85 J = 1,3
TAVE = TAVE + V(I,J)*T(I,J) /VTOT
85 CONTINUE
80 CONTINUE
WRITE (6,100) L, T, TAVE, TA3
IF (TAVE .LE. -5.0) GO TO 210

70 CONTINUE
210 CONTINUE
STOP
END
REFERENCES CITED


4. Ibid., pp. 27.1-27.5.

5. Ibid., p. 22.6.


